

THE IMPULSE RESPONSE FUNCTION

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Introduction

Consider a single-degree-of-freedom system.

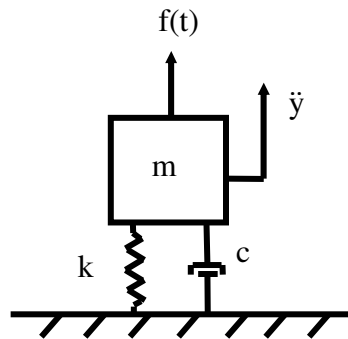


Figure 1.

The variables are

- m is the mass
- c is the viscous damping coefficient
- k is the stiffness
- y is the absolute displacement of the mass
- f(t) is the applied force

The free-body diagram is

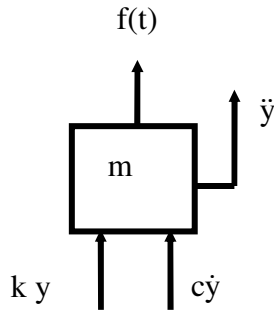


Figure 2.

Note that the double-dot denotes acceleration.

Summation of forces in the vertical direction

$$\sum F = m\ddot{y} \quad (1)$$

$$m\ddot{y} = -c\dot{y} - ky + f(t) \quad (2)$$

$$m\ddot{y} + c\dot{y} + ky = f(t) \quad (3)$$

Divide through by m ,

$$\ddot{y} + \left(\frac{c}{m}\right)\dot{y} + \left(\frac{k}{m}\right)y = \left(\frac{1}{m}\right)f(t) \quad (4)$$

By convention,

$$(c/m) = 2\xi\omega_n \quad (5)$$

$$(k/m) = \omega_n^2 \quad (6)$$

where

ω_n is the natural frequency in (radians/sec)

ξ is the damping ratio

By substitution,

$$\ddot{y} + 2\xi\omega_n \dot{y} + \omega_n^2 y = \frac{1}{m} f(t) \quad (7)$$

Now assume an impulse force using the Dirac delta function.

$$f(t) = \delta(t) \quad (8)$$

The governing equation becomes.

$$\ddot{y} + 2\xi\omega_n \dot{y} + \omega_n^2 x = \frac{1}{m} \delta(t) \quad (9)$$

The right-hand-side can be rewritten as

$$\ddot{y} + 2\xi\omega_n \dot{y} + \omega_n^2 y = \frac{\omega_n^2}{k} \delta(t) \quad (10)$$

Take the Laplace transform of each side.

$$L\{\ddot{y} + 2\xi\omega_n \dot{y} + \omega_n^2 y\} = L\left\{\frac{\omega_n^2}{k} \delta(t)\right\} \quad (11)$$

$$\begin{aligned} & s^2 Y(s) - sy(0) - y'(0) \\ & + 2\xi\omega_n s Y(s) - 2\xi\omega_n y(0) \\ & + \omega_n^2 Y(s) = \frac{\omega_n^2}{k} \end{aligned} \quad (12)$$

$$\left\{s^2 + 2\xi\omega_n s + \omega_n^2\right\} Y(s) - \{s + 2\xi\omega_n\} y(0) - y'(0) = \frac{\omega_n^2}{k} \quad (13)$$

$$\left\{s^2 + 2\xi\omega_n s + \omega_n^2\right\} Y(s) = \frac{\omega_n^2}{k} + \{s + 2\xi\omega_n\} y(0) + y'(0) \quad (14)$$

$$s^2 + 2\xi\omega_n s + \omega_n^2 = (s + \xi\omega_n)^2 - (\xi\omega_n)^2 + \omega_n^2 \quad (15)$$

$$s^2 + 2\xi\omega_n s + \omega_n^2 = (s + \xi\omega_n)^2 + \omega_n^2(1 - \xi^2) \quad (16)$$

Let

$$\omega_d = \omega_n \sqrt{1 - \xi^2} \quad (17)$$

Substitute equation (17) into (16).

$$s^2 + 2\xi\omega_n s + \omega_n^2 = (s + \xi\omega_n)^2 + \omega_d^2 \quad (18)$$

$$\left\{ (s + \xi\omega_n)^2 + \omega_d^2 \right\} Y(s) = \frac{\omega_n^2}{k} + \{s + 2\xi\omega_n\}y(0) + y'(0) \quad (19)$$

$$\begin{aligned} Y(s) = & \frac{\omega_n^2}{k} \left\{ \frac{1}{(s + \xi\omega_n)^2 + \omega_d^2} \right\} \\ & + \left\{ \frac{s + 2\xi\omega_n}{(s + \xi\omega_n)^2 + \omega_d^2} \right\} y(0) + \left\{ \frac{1}{(s + \xi\omega_n)^2 + \omega_d^2} \right\} y'(0) \end{aligned} \quad (20)$$

Assume that the initial conditions are zero.

$$Y(s) = \frac{\omega_n^2}{k} \left\{ \frac{1}{(s + \xi\omega_n)^2 + \omega_d^2} \right\} \quad (21)$$

Take the inverse Laplace transform.

$$y(t) = \frac{\omega_n^2}{k \omega_d} \exp(-\xi \omega_n t) \sin(\omega_d t) \quad (22)$$

$$y(t) = \frac{1}{m \omega_d} \exp(-\xi \omega_n t) \sin(\omega_d t) \quad (23)$$

The impulse response function $h(t)$ is

$$h(t) = \frac{1}{m \omega_d} \exp(-\xi \omega_n t) \sin(\omega_d t) \quad , \quad t > 0 \quad (24)$$

Consider a system subjected to a force which varies arbitrarily with the time. The response of the mass can be calculated using the impulse response function embedded in a convolution integral.

Reference

1. T. Irvine, The Steady-State Response of a Single-degree-of-freedom System Subjected to a Harmonic Force, Revision A, 2000.
2. D.E. Newland, An Introduction to Random Vibrations, Spectral & Wavelet Analysis, Third Edition, New York, 2005.
3. T. Irvine, Partial Fractions in Shock and Vibration Analysis, Rev F, Vibrationdata, 2010.

APPENDIX A

The impulse response function can also be found from the steady-state transfer function.

The transfer function $H(\omega)$ for the system in Figure 1 was derived in Reference 1.

$$H(\omega) = \frac{X(\omega)}{F(\omega)} = \frac{\omega_n^2}{k} \left[\frac{-1}{\omega^2 - j\omega(2\xi\omega_n) - \omega_n^2} \right] \quad (\text{A-1})$$

The denominator on the right hand side can be factored such that

$$H(\omega) = -\frac{\omega_n^2}{k} \left[\frac{1}{\left(\omega - j\xi\omega_n + \omega_d\right)\left(\omega - j\xi\omega_n - \omega_d\right)} \right] \quad (\text{A-2})$$

Partial fraction expansion yields

$$H(\omega) = \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \left\{ \frac{1}{\omega - j\xi\omega_n + \omega_d} - \frac{1}{\omega - j\xi\omega_n - \omega_d} \right\} \quad (\text{A-3})$$

Reference 2 shows that impulse response function is the inverse Fourier transform of the transfer function.

$$h(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} H(\omega) \exp(j\omega t) d\omega \quad (\text{A-4})$$

Further note that $t > 0$.

$$h(t) = \frac{1}{2\pi} \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{\omega - j\xi\omega_n + \omega_d} - \frac{1}{\omega - j\xi\omega_n - \omega_d} \right\} \exp(j\omega t) d\omega \quad (\text{A-5})$$

$$\begin{aligned} h(t) &= \frac{1}{2\pi} \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{\omega - j\xi\omega_n + \omega_d} \right\} \exp(j\omega t) d\omega \\ &\quad - \frac{1}{2\pi} \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{\omega - j\xi\omega_n - \omega_d} \right\} \exp(j\omega t) d\omega \end{aligned} \quad (\text{A-6})$$

Substitute variables as follows

$$u = \omega - j\xi\omega_n + \omega_d \quad (\text{A-7})$$

$$\omega = u + j\xi\omega_n - \omega_d \quad (\text{A-8})$$

$$du = d\omega \quad (\text{A-9})$$

$$v = \omega - j\xi\omega_n - \omega_d \quad (\text{A-10})$$

$$\omega = v + j\xi\omega_n + \omega_d \quad (\text{A-11})$$

$$dv = d\omega \quad (\text{A-12})$$

$$\begin{aligned}
h(t) = & \frac{1}{2\pi} \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{u} \right\} \exp\left(j \left[u + j\xi\omega_n - \omega_d \right] t \right) du \\
& - \frac{1}{2\pi} \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{v} \right\} \exp\left(j \left[v + j\xi\omega_n + \omega_d \right] t \right) dv
\end{aligned}
\tag{A-13}$$

$$\begin{aligned}
h(t) = & \frac{1}{2\pi} \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{u} \right\} \exp\left(\left[ju - \xi\omega_n - j\omega_d \right] t \right) du \\
& - \frac{1}{2\pi} \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{v} \right\} \exp\left(\left[jv - \xi\omega_n + j\omega_d \right] t \right) dv
\end{aligned}
\tag{A-14}$$

$$\begin{aligned}
h(t) = & \frac{1}{2\pi} \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{u} \right\} \exp(jut) \exp\left(-\xi\omega_n t \right) \exp\left(-j\omega_d t \right) du \\
& - \frac{1}{2\pi} \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{v} \right\} \exp(jvt) \exp\left(-\xi\omega_n t \right) \exp\left(+j\omega_d t \right) dv
\end{aligned}
\tag{A-15}$$

$$\begin{aligned}
h(t) = & \frac{1}{2\pi} \exp(-\xi\omega_n t) \exp(-j\omega_d t) \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{u} \right\} \exp(jut) du \\
& - \frac{1}{2\pi} \exp(-\xi\omega_n t) \exp(+j\omega_d t) \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{v} \right\} \exp(jvt) dv
\end{aligned} \tag{A-16}$$

Given that limits go to infinity, the two integral terms are the same.

$$\begin{aligned}
h(t) = & \frac{1}{2\pi} \exp(-\xi\omega_n t) \exp(-j\omega_d t) \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{u} \right\} \exp(jut) du \\
& - \frac{1}{2\pi} \exp(-\xi\omega_n t) \exp(+j\omega_d t) \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{u} \right\} \exp(jut) du
\end{aligned} \tag{A-17}$$

$$h(t) = \frac{1}{2\pi} \exp(-\xi\omega_n t) \left\{ \exp(-j\omega_d t) - \exp(+j\omega_d t) \right\} \left\{ \frac{\omega_n^2}{2k\omega_d} \right\} \int_{-\infty}^{\infty} \left\{ \frac{1}{u} \right\} \exp(jut) du \tag{A-18}$$

$$h(t) = -j \frac{1}{\pi} \left\{ \frac{\omega_n^2}{k\omega_d} \right\} \exp(-\xi\omega_n t) \sin(\omega_d t) \int_{-\infty}^{\infty} \left\{ \frac{1}{u} \right\} \exp(jut) du \tag{A-19}$$

$$h(t) = -j \frac{1}{\pi} \left\{ \frac{1}{m \omega_d} \right\} \exp(-\xi \omega_n t) \sin(\omega_d t) \int_{-\infty}^{\infty} \left\{ \frac{1}{u} \right\} \exp(jut) du \quad (\text{A-20})$$

$$h(t) = -j \frac{1}{\pi} \left\{ \frac{1}{m \omega_d} \right\} \exp(-\xi \omega_n t) \sin(\omega_d t) \int_{-\infty}^{\infty} \left\{ \frac{1}{u} \right\} \{\cos(ut) + j \sin(ut)\} du \quad (\text{A-21})$$

Again note that $t > 0$.

$$h(t) = -j \frac{2}{\pi} \left\{ \frac{1}{m \omega_d} \right\} \exp(-\xi \omega_n t) \sin(\omega_d t) \int_{0^+}^{\infty} \left\{ \frac{1}{u} \right\} \{\cos(ut) + j \sin(ut)\} du \quad (\text{A-22})$$

$$h(t) = \frac{2}{\pi} \left\{ \frac{1}{m \omega_d} \right\} \exp(-\xi \omega_n t) \sin(\omega_d t) \int_{0^+}^{\infty} \left\{ \frac{1}{u} \right\} \{-j \cos(ut) + \sin(ut)\} du \quad (\text{A-23})$$

The impulse response function must be real.

$$h(t) = \frac{2}{\pi} \left\{ \frac{1}{m \omega_d} \right\} \exp(-\xi \omega_n t) \sin(\omega_d t) \int_{0^+}^{\infty} \left\{ \frac{1}{u} \sin(ut) \right\} du \quad (\text{A-24})$$

From integral tables,

$$\int_{0^+}^{\infty} \left\{ \frac{1}{u} \sin(ut) \right\} du = \frac{\pi}{2} \quad (\text{A-25})$$

The impulse response function is thus

$$h(t) = \left\{ \frac{1}{m \omega_d} \right\} \exp(-\xi \omega_n t) \sin(\omega_d t) \quad (\text{A-26})$$