## DERIVATION OF MILES EQUATION FOR THE RELATIVE DISPLACEMENT RESPONSE TO BASE EXCITATION

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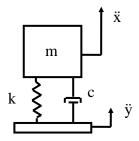
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## Introduction

The objective is to derive a Miles equation which gives the overall relative displacement response of a single-degree-of-freedom system to an applied force, where the excitation is in the form of a random vibration acceleration power spectral density.

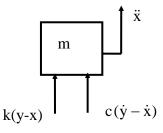
## Derivation

Consider a single degree-of-freedom system



where m equals mass, c equals the viscous damping coefficient, and k equals the stiffness. The absolute displacement of the mass equals x, and the base input displacement equals y.

The free-body diagram is



Summation of forces in the vertical direction,

$$\sum F = m\ddot{x} \tag{1}$$

$$m\ddot{\mathbf{x}} = \mathbf{c}(\dot{\mathbf{y}} - \dot{\mathbf{x}}) + \mathbf{k}(\mathbf{y} - \mathbf{x}) \tag{2}$$

Substituting the relative displacement terms into equation (2) yields

$$m(\ddot{z} + \ddot{y}) = -c\dot{z} - kz \tag{3}$$

$$m\ddot{z} + c\dot{z} + kz = -m\ddot{y} \tag{4}$$

Dividing through by mass yields,

$$\ddot{z} + (c/m)\dot{z} + (k/m)z = -\ddot{y} \tag{5}$$

By convention,

$$(c/m) = 2\xi\omega_{n}$$
$$(k/m) = \omega_{n}^{2}$$

where  $\omega_n$  is the natural frequency in (radians/sec), and  $\xi$  is the damping ratio.

$$\ddot{z} + 2\xi \omega_n \, \dot{z} + \omega_n^2 z = -\ddot{y} \tag{6}$$

Let

$$z(t) = \hat{Z}\exp(j\omega t) \tag{7}$$

$$y(t) = \hat{Y} \exp(j\omega t)$$
 (8)

$$\left[-\omega^2 + j2\xi\omega\omega_n + \omega_n^2\right] \hat{Z}\exp(j\omega t) = -\hat{Y}\exp(j\omega t)$$
 (9)

Take the Fourier transform of each side.

$$\left[-\omega^2 + j2\xi\omega\omega_n + \omega_n^2\right] Z(\omega) = -Y(\omega)$$
 (10)

$$Z(\omega) = \frac{-Y(\omega)}{\left[-\omega^2 + \omega_n^2 + j2\xi\omega\omega_n\right]}$$
(11)

$$Z(f) = \frac{-(1/2\pi)Y(f)}{\left[-f^2 + f_n^2 + j2\xi f f_n\right]}$$
(12)

Divide through by  $f_n^2$ .

$$Z(f) = \frac{-\left(1/(2\pi f_n)^2\right)Y(f)}{\left[-\rho^2 + 1 + j2\xi\rho\right]}$$
(13)

where  $\rho = f/f_n$ 

$$Z(f) = \frac{-\left(1/(2\pi f_n)^2\right)Y(f)}{\left[1-\rho^2 + j2\xi\rho\right]}$$
(14)

Multiply each side by its complex conjugate.

$$Z(f)Z^*(f) = \frac{\left(1/(2\pi f_n)^2\right)}{\left[1-\rho^2 + j2\xi\rho\right]} \frac{\left(1/(2\pi f_n)^2\right)}{\left[1-\rho^2 - j2\xi\rho\right]} Y(f)Y^*(f)$$
(15)

$$Z(f)Z*(f) = \frac{\left(1/(2\pi f_n)^2\right)}{\left[\rho^2 - 1 + j2\xi\rho\right]} \frac{\left(1/(2\pi f_n)^2\right)}{\left[\rho^2 - 1 - j2\xi\rho\right]} Y(f)Y*(f)$$
(16)

The transfer function  $H(\rho)$  times its complex conjugate is

$$H(\rho)H*(\rho) = \left[\frac{\left(1/(2\pi f_n)^2\right)}{\rho^2 - j2\xi\rho - 1}\right] \left[\frac{\left(1/(2\pi f_n)^2\right)}{\rho^2 + j2\xi\rho - 1}\right]$$
(17)

Solve for the roots R1 and R2 of the first denominator.

$$R1, R2 = \frac{j2\xi \pm \sqrt{(-j2\xi)^2 - 4(-1)}}{2}$$
 (18)

$$R1, R2 = \frac{j2\xi \pm \sqrt{-4\xi^2 + 4}}{2} \tag{19}$$

$$R1, R2 = j\xi \pm \sqrt{1 - \xi^2}$$
 (20)

Solve for the roots R3 and R4 of the second denominator.

$$R3, R4 = \frac{-j\xi \pm \sqrt{(j2\xi)^2 - 4(-1)}}{2}$$
 (21)

$$R3, R4 = \frac{-j2\xi \pm \sqrt{-4\xi^2 + 4}}{2}$$
 (22)

$$R3, R4 = -j\xi \pm \sqrt{1 - \xi^2}$$
 (23)

(24)

Summary,

$$R1 = +j\xi + \sqrt{1-\xi^2}$$
 (25)

$$R2 = +j\xi - \sqrt{1 - \xi^2}$$
 (26)

$$R3 = -j\xi + \sqrt{1 - \xi^2}$$
 (27)

$$R4 = -j\xi - \sqrt{1 - \xi^2}$$
 (28)

Note

$$R2 = -R1*$$
 (29)

$$R3 = R1*$$
 (30)

$$R4 = -R1* \tag{31}$$

Now substitute into the denominators.

$$H(\rho)H*(\rho) = \left[\frac{1/(2\pi f_n)^4}{\left(\rho - j\xi - \sqrt{1 - \xi^2}\right)\left(\rho - j\xi + \sqrt{1 - \xi^2}\right)\left(\rho + j\xi - \sqrt{1 - \xi^2}\right)\left(\rho + j\xi + \sqrt{1 - \xi^2}\right)}\right]$$
(32)

$$H(\rho)H*(\rho) = \left[\frac{1/(2\pi f_n)^4}{(\rho - R1)(\rho - R2)(\rho - R3)(\rho - R4)}\right]$$
(33)

$$H(\rho)H*(\rho) = \left[\frac{1/(2\pi f_n)^4}{(\rho - R1)(\rho + R1*)(\rho - R1*)(\rho + R1)}\right]$$
(34)

Expand into partial fractions.

$$\left[\frac{1}{(\rho - R1)(\rho + R1^*)(\rho - R1^*)(\rho + R1)}\right] = +\frac{\alpha}{(\rho - R1)} + \frac{\beta}{(\rho - R1^*)} + \frac{\lambda}{(\rho + R1^*)} + \frac{\sigma}{(\rho + R1)}$$

(35)

Equation (35) is solved using the method in Reference 1.

$$\begin{bmatrix} \alpha \\ \beta \\ \lambda \\ \sigma \end{bmatrix} = \frac{1}{\left[8\xi\sqrt{1-\xi^{2}}\right]} \begin{bmatrix} -\xi - j\sqrt{1-\xi^{2}} \\ -\xi + j\sqrt{1-\xi^{2}} \\ +\xi - j\sqrt{1-\xi^{2}} \\ +\xi + j\sqrt{1-\xi^{2}} \end{bmatrix}$$

$$(36)$$

Let

$$\psi = \frac{\xi}{\sqrt{1 - \xi^2}} \tag{37}$$

$$\begin{bmatrix} \alpha \\ \beta \\ \lambda \\ \sigma \end{bmatrix} = \frac{1}{8\xi} \begin{bmatrix} -\psi - j \\ -\psi + j \\ +\psi - j \\ +\psi + j \end{bmatrix}$$
(38)

$$R1 = +j\xi + \sqrt{1 - \xi^2}$$
 (39)

$$\left[\frac{1}{(\rho - R1)(\rho + R1^*)(\rho - R1^*)(\rho + R1)}\right] = + \frac{-\psi - j}{\left(\rho - \sqrt{1 - \xi^2} - j\xi\right)} \left[\frac{1}{8\xi}\right] + \frac{-\psi + j}{\left(\rho - \sqrt{1 - \xi^2} + j\xi\right)} \left[\frac{1}{8\xi}\right] + \frac{+\psi - j}{\left(\rho + \sqrt{1 - \xi^2} - j\xi\right)} \left[\frac{1}{8\xi}\right] + \frac{+\psi + j}{\left(\rho + \sqrt{1 - \xi^2} + j\xi\right)} \left[\frac{1}{8\xi}\right]$$

$$+ \frac{+\psi + j}{\left(\rho + \sqrt{1 - \xi^2} + j\xi\right)} \left[\frac{1}{8\xi}\right]$$

$$(40)$$

The overall relative displacement is found by integration.

$$\left[z_{RMS}(f_n,\xi)\right]^2 = \int_0^\infty H(\rho)H^*(\rho) \,\hat{Y}_{APSD}(f) df \tag{41}$$

$$\left[z_{RMS}(f_n,\xi)\right]^2 = f_n \int_0^\infty H(\rho)H^*(\rho) \,\hat{Y}_{APSD}(f) \,d\rho \tag{42}$$

Assume that the acceleration PSD is constant

$$\hat{Y}_{APSD}(f) = A \tag{43}$$

$$\left[x_{RMS}(f_n,\xi)\right]^2 = A f_n \int_0^\infty H(\rho) H^*(\rho) d\rho \tag{44}$$

$$\left\{ x_{RMS} \left( f_{n}, \xi \right) \right\}^{2} \left[ \frac{8\xi (2\pi)^{4} f_{n}^{3}}{A} \right] = + \int_{0}^{\infty} \frac{-\psi - j}{\left( \rho - \sqrt{1 - \xi^{2}} - j\xi \right)} d\rho$$

$$+ \int_{0}^{\infty} \frac{-\psi + j}{\left( \rho - \sqrt{1 - \xi^{2}} + j\xi \right)} d\rho$$

$$+ \int_{0}^{\infty} \frac{+\psi - j}{\left( \rho + \sqrt{1 - \xi^{2}} - j\xi \right)} d\rho$$

$$+ \int_{0}^{\infty} \frac{+\psi + j}{\left( \rho + \sqrt{1 - \xi^{2}} + j\xi \right)} d\rho$$

$$(45)$$

Equation (45) is solved using the method in Reference 1.

$$\left\{ z_{\text{RMS}} \left( f_{\text{n}}, \xi \right) \right\}^{2} \left[ \frac{8\xi (2\pi)^{4} f_{\text{n}}^{3}}{A} \right] = 2\pi$$
 (46)

$$\left\{ z_{RMS} \left( f_{\mathbf{n}}, \xi \right) \right\}^{2} = \left[ \frac{\pi}{4\xi} \right] \frac{A}{\left( 2\pi \right)^{4} f_{\mathbf{n}}^{3}}$$

$$\tag{47}$$

$$z_{\text{RMS}}(f_n, \xi) = \sqrt{\frac{\pi A}{4\xi(2\pi)^4 f_n^3}}$$
 (48)

The RMS relative displacement is finally

$$z_{RMS}(f_n, \xi) = \sqrt{\frac{A}{64\pi^3 \xi f_n^3}}$$
 (49)

$$Q = \frac{1}{2\xi} \tag{50}$$

$$z_{RMS} = \sqrt{\frac{QA}{32\pi^3 f_n^3}}$$
 (51)

Note that equation for the absolute acceleration response in Reference 2 is

$$\ddot{\mathbf{x}}_{\mathbf{GRMS}}(\mathbf{f}_{\mathbf{n}}, \mathbf{Q}) = \sqrt{\left(\frac{\pi}{2}\right)} \mathbf{f}_{\mathbf{n}} \mathbf{Q} \mathbf{A} \tag{52}$$

Thus

$$\ddot{\mathbf{x}}_{\mathbf{GRMS}} = [2\pi \mathbf{f}_{\mathbf{n}}]^2 \mathbf{z}_{\mathbf{RMS}} \tag{53}$$

## References

- 1. T. Irvine, Derivation of Miles Equation for an Applied Force, Vibrationdata, 2008.
- 2. T. Irvine, Derivation of Miles Equation Revision D, Vibrationdata, 2008.