

# LONGITUDINAL NATURAL FREQUENCIES OF RODS AND RESPONSE TO INITIAL CONDITIONS

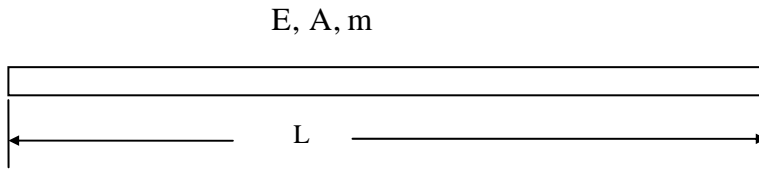
Revision B

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Consider a thin rod.



$E$  is the modulus of elasticity.

$A$  is the cross-section area.

$m$  is the mass per unit length.

The longitudinal displacement  $u(x, t)$  is governed by the equation

$$\frac{\partial}{\partial x} \left[ EA(x) \frac{\partial u}{\partial x} \right] = m(x) \frac{\partial^2 u}{\partial t^2} \quad (1)$$

This equation is taken from Reference 1.

For a uniform cross-section, the governing equation simplifies to

$$EA(x) \frac{\partial^2 u}{\partial x^2} = m(x) \frac{\partial^2 u}{\partial t^2} \quad (2)$$

Consider a beam with uniform mass density. The governing equation simplifies to

$$\frac{\partial^2 u}{\partial x^2} = \left( \frac{\rho}{E} \right) \frac{\partial^2 u}{\partial t^2} \quad (3)$$

where

$\rho$  is the mass per unit volume.

Let

$$c = \sqrt{\frac{E}{\rho}} \quad (4)$$

Note that  $c$  is the longitudinal wave velocity. Substitute equation (4) into (3).

$$\frac{\partial^2 u}{\partial x^2} = \left( \frac{1}{c^2} \right) \frac{\partial^2 u}{\partial t^2} \quad (5)$$

Separate the variables. Let

$$u(x, t) = U(x)T(t) \quad (6)$$

Substitute equation (6) into (5).

$$\frac{\partial^2}{\partial x^2} [U(x)T(t)] = \left( \frac{1}{c^2} \right) \frac{\partial^2}{\partial t^2} [U(x)T(t)] \quad (7)$$

Perform the partial differentiation.

$$U''(x)T(t) = \left( \frac{1}{c^2} \right) U(x)T''(t) \quad (8)$$

Divide through by  $U(x)T(t)$ .

$$\frac{U''(x)}{U(x)} = \left( \frac{1}{c^2} \right) \frac{T''(t)}{T(t)} \quad (9)$$

$$c^2 \frac{U''(x)}{U(x)} = \frac{T''(t)}{T(t)} \quad (10)$$

Each side of equation (10) must equal a constant. Let  $\omega$  be a constant.

$$c^2 \frac{U''(x)}{U(x)} = \frac{T''(t)}{T(t)} = -\omega^2 \quad (11)$$

The time equation is

$$\frac{T''(t)}{T(t)} = -\omega^2 \quad (12)$$

$$T''(t) = -\omega^2 T(t) \quad (13)$$

$$T''(t) + \omega^2 T(t) = 0 \quad (14)$$

Propose a solution

$$T(t) = a \sin(\omega t) + b \cos(\omega t) \quad (15)$$

$$T'(t) = a \omega \cos(\omega t) - b \omega \sin(\omega t) \quad (16)$$

$$T''(t) = -a \omega^2 \sin(\omega t) - b \omega^2 \cos(\omega t) \quad (17)$$

Verify the proposed solution. Substitute into equation (14).

$$-a \omega^2 \sin(\omega t) - b \omega^2 \cos(\omega t) + \omega^2 [\sin(\omega t) + \omega^2 \cos(\omega t)] = 0 \quad (18)$$

$$0 = 0 \quad (19)$$

Equation (15) is thus verified as a solution.

There is not a unique  $\omega$ , however, in equation (11). This is demonstrated later in the derivation. Thus a subscript  $n$  must be added as follows.

$$T_n(t) = a_n \sin(\omega_n t) + b_n \cos(\omega_n t) \quad (20)$$

The spatial equation is

$$c^2 \frac{U''(x)}{U(x)} = -\omega^2 \quad (21)$$

$$c^2 U''(x) = -\omega^2 U(x) \quad (22)$$

$$c^2 U''(x) + \omega^2 U(x) = 0 \quad (23)$$

$$U''(x) + \frac{\omega^2}{c^2} U(x) = 0 \quad (24)$$

Equation (24) is similar to equation (14). Thus, a solution can be found by inspection.

$$U(x) = d \sin\left(\frac{\omega x}{c}\right) + e \cos\left(\frac{\omega x}{c}\right) \quad (25)$$

The slope equation is

$$U'(x) = \left[\frac{\omega}{c}\right] \left[ d \cos\left(\frac{\omega x}{c}\right) - e \sin\left(\frac{\omega x}{c}\right) \right] \quad (26)$$

Now consider three boundary condition cases as shown in the following appendices.

#### REFERENCE

1. L. Meirovitch, Analytical Methods in Vibrations, Macmillan, New York, 1967.

## APPENDIX A

### Case I. Fixed-Fixed

The left boundary condition is

$$u(0, t) = 0 \quad (\text{zero displacement}) \quad (\text{A-1})$$

$$U(0)T(t) = 0 \quad (\text{A-2})$$

$$U(0) = 0 \quad (\text{A-3})$$

The right boundary condition is

$$u(L, t) = 0 \quad (\text{zero displacement}) \quad (\text{A-4})$$

$$U(L)T(t) = 0 \quad (\text{A-5})$$

$$U(L) = 0 \quad (\text{A-6})$$

Substitute equation (A-3) into (25).

$$e = 0 \quad (\text{A-7})$$

Thus, the displacement equation becomes

$$U(x) = d \sin\left(\frac{\omega x}{c}\right) \quad (\text{A-8})$$

Substitute equation (A-6) into (A-8).

$$d \sin\left(\frac{\omega L}{c}\right) = 0 \quad (\text{A-9})$$

The constant  $d$  must be non-zero for a non-trivial solution. Thus,

$$\frac{\omega_n L}{c} = n\pi, \quad n = 1, 2, 3, \dots \quad (\text{A-10})$$

The  $\omega$  term is given a subscript  $n$  because there are multiple roots.

$$\omega_n = n\pi \frac{c}{L}, \quad n = 1, 2, 3, \dots \quad (\text{A-11})$$

The displacement function the fixed-fixed rod is

$$U_n(x) = d_n \sin\left(\frac{\omega_n x}{c}\right) \quad (\text{A-12})$$

$$U_n(x) = d_n \sin\left(\frac{n\pi x}{L}\right) \quad (\text{A-13})$$

Substitute the natural frequency term into the time equation.

$$T_n(t) = a_n \sin\left(\frac{n\pi c t}{L}\right) + b_n \cos\left(\frac{n\pi c t}{L}\right) \quad (\text{A-14})$$

The displacement function is thus

$$u(x, t) = \sum_{n=1}^{\infty} \left\{ \left[ d_n \sin\left(\frac{n\pi x}{L}\right) \right] \left[ a_n \sin\left(\frac{n\pi c t}{L}\right) + b_n \cos\left(\frac{n\pi c t}{L}\right) \right] \right\} \quad (\text{A-15})$$

The coefficients can be simplified as follows

$$A_n = d_n a_n \quad (\text{A-16})$$

$$B_n = d_n b_n \quad (\text{A-17})$$

By substitution, the displacement equation is

$$u(x, t) = \sum_{n=1}^{\infty} \left\{ \left[ \sin\left(\frac{n\pi x}{L}\right) \right] \left[ A_n \sin\left(\frac{n\pi c t}{L}\right) + B_n \cos\left(\frac{n\pi c t}{L}\right) \right] \right\} \quad (\text{A-18})$$

## APPENDIX B

### Case II. Fixed-Free

The left boundary conditions is

$$u(0, t) = 0 \quad (\text{zero displacement}) \quad (\text{B-1})$$

$$U(0)T(t) = 0 \quad (\text{B-2})$$

$$U(0) = 0 \quad (\text{B-3})$$

The right boundary condition is

$$\left. \frac{\partial}{\partial x} u(x, t) \right|_{x=L} = 0 \quad (\text{zero stress}) \quad (\text{B-4})$$

$$U'(L)T(t) = 0 \quad (\text{B-5})$$

$$U'(L) = 0 \quad (\text{B-6})$$

Substitute equation (B-3) into (25).

$$e = 0 \quad (\text{B-7})$$

Thus, the displacement equation becomes

$$U(x) = d \sin\left(\frac{\omega x}{c}\right) \quad (\text{B-8})$$

$$U'(x) = \left[ \frac{\omega}{c} \right] \left[ d \cos\left(\frac{\omega x}{c}\right) \right] \quad (\text{B-9})$$

Substitute equation (B-6) into equation (B-9).

$$d \cos\left(\frac{\omega L}{c}\right) = 0 \quad (\text{B-10})$$

The constant  $d$  must be non-zero for a non-trivial solution. Thus,

$$\frac{\omega_n L}{c} = \left(\frac{2n-1}{2}\right)\pi, \quad n = 1, 2, 3, \dots \quad (\text{B-11})$$

The  $\omega$  term is given a subscript  $n$  because there are multiple roots.

$$\omega_n = \left(\frac{2n-1}{2}\right)\pi \frac{c}{L}, \quad n = 1, 2, 3, \dots \quad (\text{B-12})$$

The displacement function for the fixed-free rod is

$$U_n(x) = d_n \sin\left(\frac{\omega_n x}{c}\right) \quad (\text{B-13})$$

$$U_n(x) = d_n \sin\left(\frac{(2n-1)\pi x}{2L}\right) \quad (\text{B-14})$$

Substitute the natural frequency term into the time equation.

$$T_n(t) = a_n \sin\left(\frac{(2n-1)\pi c t}{2L}\right) + b_n \cos\left(\frac{(2n-1)\pi c t}{2L}\right) \quad (\text{B-15})$$

The displacement function is thus

$$u(x,t) = \sum_{n=1}^{\infty} \left\{ \left[ d_n \sin\left(\frac{(2n-1)\pi x}{2L}\right) \right] \left[ a_n \sin\left(\frac{(2n-1)\pi c t}{2L}\right) + b_n \cos\left(\frac{(2n-1)\pi c t}{2L}\right) \right] \right\} \quad (\text{B-16})$$



Simplify the coefficients.

$$u(x,t) = \sum_{n=1}^{\infty} \left\{ \left[ \sin\left(\frac{(2n-1)\pi x}{2L}\right) \right] \left[ A_n \sin\left(\frac{(2n-1)\pi c t}{2L}\right) + B_n \cos\left(\frac{(2n-1)\pi c t}{2L}\right) \right] \right\} \quad (\text{B-17})$$

Now determine the effective mass of the rod for the fundamental mode. The stiffness  $k$  at free end of the fixed-free longitudinal rod is

$$k = \frac{EA}{L} \quad (\text{B-18})$$

The formula for the fundamental frequency of a single-degree-of-freedom system is

$$\omega_1 = \sqrt{\frac{k}{m}} \quad (\text{B-19})$$

Solve for the mass  $m$ .

$$m = \frac{k}{\omega_1^2} \quad (\text{B-20})$$

Substitute the stiffness term from equation (B-18).

$$m = \frac{EA}{\omega_1^2 L} \quad (\text{B-21})$$

Add a subscript  $e$  to denote that the mass is the effective mass.

$$m_e = \frac{EA}{\omega_1^2 L} \quad (\text{B-22})$$

Calculate the fundamental frequency from equation (B-12).

$$\omega_1 = \left(\frac{1}{2}\right) \pi \frac{c}{L} \quad (\text{B-23})$$

$$\omega_1 = \left(\frac{1}{2}\right) \pi \frac{1}{L} \sqrt{\frac{E}{\rho}} \quad (\text{B-24})$$

$$\omega_1^2 = \left( \frac{\pi^2}{4L^2} \right) \frac{E}{\rho} \quad (\text{B-25})$$

Substitute the frequency term from equation (B-25) into (B-22).

$$m_e = \frac{4L^2 \rho EA}{E \pi^2 L} \quad (\text{B-26})$$

$$m_e = \frac{4L \rho A}{\pi^2} \quad (\text{B-27})$$

Let M be the mass of the rod. The effective mass is

$$m_e = \frac{4}{\pi^2} M \quad (\text{B-28})$$

## APPENDIX C

### Case III. Free-Free

The left boundary conditions is

$$\left. \frac{\partial}{\partial x} u(x, t) \right|_{x=0} = 0 \quad (\text{zero stress}) \quad (\text{C-1})$$

$$U'(0)T(t) = 0 \quad (\text{C-2})$$

$$U'(0) = 0 \quad (\text{C-3})$$

The right boundary condition is

$$\left. \frac{\partial}{\partial x} u(x, t) \right|_{x=L} = 0 \quad (\text{zero stress}) \quad (\text{C-4})$$

$$U'(L)T(t) = 0 \quad (\text{C-5})$$

$$U'(L) = 0 \quad (\text{C-6})$$

Apply equation (C-3) to (25).

$$d = 0 \quad (\text{C-7})$$

Thus

$$U(x) = \cos\left(\frac{\omega x}{c}\right) \quad (\text{C-8})$$

The slope equation is

$$U'(x) = -\left[\frac{\omega}{c}\right] \left[ e \sin\left(\frac{\omega x}{c}\right) \right] \quad (\text{C-9})$$

Substitute equation (C-6) into (C-9).

$$e \sin\left(\frac{\omega L}{c}\right) = 0 \quad (\text{C-10})$$

The constant  $e$  must be non-zero for a non-trivial solution. Thus,

$$\frac{\omega_n L}{c} = n\pi, \quad n = 1, 2, 3, \dots \quad (\text{C-11})$$

The  $\omega$  term is given a subscript  $n$  because there are multiple roots.

$$\omega_n = n\pi \frac{c}{L}, \quad n = 1, 2, 3, \dots \quad (\text{C-12})$$

The displacement function for the free-free rod is

$$U_n(x) = e_n \cos\left(\frac{\omega_n x}{c}\right) \quad (\text{C-13})$$

$$U_n(x) = e_n \cos\left(\frac{n\pi x}{L}\right) \quad (\text{C-14})$$

Substitute the natural frequency term into the time equation.

$$T_n(t) = a_n \sin\left(\frac{n\pi c t}{L}\right) + b_n \cos\left(\frac{n\pi c t}{L}\right) \quad (\text{C-15})$$

The displacement function is thus

$$u(x, t) = \sum_{n=1}^{\infty} \left\{ \left[ e_n \cos\left(\frac{n\pi x}{L}\right) \right] \left[ a_n \sin\left(\frac{n\pi c t}{L}\right) + b_n \cos\left(\frac{n\pi c t}{L}\right) \right] \right\} \quad (\text{C-16})$$

Simplify the coefficients.

$$u(x, t) = \sum_{n=1}^{\infty} \left\{ \left[ \cos\left(\frac{n\pi x}{L}\right) \right] \left[ A_n \sin\left(\frac{n\pi c t}{L}\right) + B_n \cos\left(\frac{n\pi c t}{L}\right) \right] \right\} \quad (\text{C-17})$$

## APPENDIX D

### Fixed-Free Rod Subjected to Initial Displacement and Initial Velocity

Recall

$$u(x, t) = \sum_{n=1}^{\infty} \left\{ \sin\left(\frac{\omega_n x}{c}\right) [A_n \sin(\omega_n t) + B_n \cos(\omega_n t)] \right\} \quad (D-1)$$

Let

$$u(x, 0) = f(x) \quad (D-2)$$

$$\dot{u}(x, 0) = g(x) \quad (D-3)$$

$$f(x) = \sum_{n=1}^{\infty} \left\{ B_n \sin\left(\frac{\omega_n x}{c}\right) \right\} \quad (D-4)$$

Premultiply by  $\sin\left(\frac{\omega_m x}{c}\right)$  and integrate.

$$\int_0^L f(x) \sin\left(\frac{\omega_m x}{c}\right) dx = \sum_{n=1}^{\infty} \left\{ B_n \int_0^L \sin\left(\frac{\omega_m x}{c}\right) \sin\left(\frac{\omega_n x}{c}\right) dx \right\} \quad (D-5)$$

$$\omega_m = \left( \frac{2m-1}{2} \right) \pi \frac{c}{L}, \quad m = 1, 2, 3, \dots \quad (D-6)$$

For  $m \neq n$ ,

The integral on the right hand side of (D-5) goes to zero. The steps are omitted for brevity.

For  $m = n$ ,

$$\int_0^L f(x) \sin\left(\frac{\omega_m x}{c}\right) dx = B_m \int_0^L \sin^2\left(\frac{\omega_m x}{c}\right) dx \quad (D-7)$$

$$\int_0^L f(x) \sin\left(\frac{\omega_m x}{c}\right) dx = \frac{1}{2} B_m \int_0^L \left[1 - \cos\left(\frac{2\omega_m x}{c}\right)\right] dx \quad (D-8)$$

$$\int_0^L f(x) \sin\left(\frac{\omega_m x}{c}\right) dx = \frac{1}{2} B_m \left[ x - \left(\frac{c}{2\omega_m}\right) \sin\left(\frac{2\omega_m x}{c}\right) \right]_0^L \quad (D-9)$$

$$\int_0^L f(x) \sin\left(\frac{\omega_m x}{c}\right) dx = \frac{1}{2} B_m \left[ L - \left(\frac{c}{2\omega_m}\right) \sin\left(\frac{2\omega_m L}{c}\right) \right] \quad (D-10)$$

$$\omega_m = \left(\frac{2m-1}{2}\right) \pi \frac{c}{L}, \quad m = 1, 2, 3, \dots \quad (D-11)$$

$$\int_0^L f(x) \sin\left(\frac{\omega_m x}{c}\right) dx = \frac{1}{2} B_m L \quad (D-12)$$

$$B_m = \frac{2}{L} \int_0^L f(x) \sin\left(\frac{\omega_m x}{c}\right) dx \quad (D-13)$$

$$\dot{u}(x, t) = \sum_{n=1}^{\infty} \left\{ \omega_n \sin\left(\frac{\omega_n x}{c}\right) [A_n \cos(\omega_n t) - B_n \sin(\omega_n t)] \right\} \quad (D-14)$$

$$g(x) = \sum_{n=1}^{\infty} \left\{ \omega_n \sin\left(\frac{\omega_n x}{c}\right) [A_n] \right\} \quad (D-15)$$

$$\int_0^L g(x) \sin\left(\frac{\omega_m x}{c}\right) dx = \omega_m \sum_{n=1}^{\infty} \left\{ A_n \int_0^L \sin\left(\frac{\omega_m x}{c}\right) \sin\left(\frac{\omega_n x}{c}\right) dx \right\} \quad (D-16)$$

Equation (D-16) is similar to (D-5).

$$A_m = \frac{2}{\omega_m L} \int_0^L f(x) \sin\left(\frac{\omega_m x}{c}\right) dx \quad (D-17)$$

Again,

$$\omega_m = \left( \frac{2m-1}{2} \right) \pi \frac{c}{L}, \quad m = 1, 2, 3, \dots \quad (D-11)$$