# VIBRATION ANALYSIS OF A MASS WITH AN ARBITRARY NUMBER OF ISOLATORS <br> Revision B 

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## Introduction

An avionics component may be mounted with isolator grommets, which act as soft springs. The goal of the isolator design is to provide attenuation of shock and vibration energy. This is achieved primarily by lowering the natural frequency of the component system. In addition, the isolators provide damping.

Consider a component with a complex geometry that is to be mounted via an arbitrary number of isolators, as shown in Figures 1 and 2. Assume that the component's hardmounted natural frequency is at least one octave greater than any of its isolation frequencies.

The objective is to derive the equations of motion for this system, accounting for six degrees-of-freedom.

## Derivation



Figure 1. Isolated Avionics Component Model
The mass and inertia are represented at a point with the circle symbol. Each isolator is modeled by three orthogonal DOF springs. Only one spring set is shown for brevity. The triangles indicate fixed constraints. The origin is at the center of gravity.
$\underset{\mathrm{z}}{\substack{\mathrm{y}}} \mathrm{x}$


$$
k_{y, i}\left(-y+c_{i} \alpha-a_{i} \theta\right)
$$

Figure 2. Isolated Avionics Component Model with Dimensions

The variables $\alpha, \beta$, and $\theta$ represent rotations about the $\mathrm{X}, \mathrm{Y}$, and Z axes, respectively, using the right-hand rule convention.

Derive the equations of motion. Let n be the number of isolators.

$$
\begin{align*}
& m \ddot{x}=\sum_{i=1}^{n} k_{x, i}\left(-x-c_{i} \beta-b_{i} \theta\right)  \tag{1}\\
& m \ddot{x}+\sum_{i=1}^{n} k_{x, i}\left(x+c_{i} \beta+b_{i} \theta\right)=0  \tag{2}\\
& m \ddot{x}+x \sum_{i=1}^{n} k_{x, i}+\beta \sum_{i=1}^{n} k_{x, i} c_{i}+\theta \sum_{i=1}^{n} k_{x, i} b_{i}=0  \tag{3}\\
& m \ddot{y}=\sum_{i=1}^{n} k_{y, i}\left(-y+c_{i} \alpha-a_{i} \theta\right)  \tag{4}\\
& m \ddot{y}+\sum_{i=1}^{n} k_{y, i}\left(y-c_{i} \alpha+a_{i} \theta\right)=0  \tag{5}\\
& m \ddot{y}+y \sum_{i=1}^{n} k_{y, i}-\alpha \sum_{i=1}^{n} k_{y, i} c_{i}+\theta \sum_{i=1}^{n} k_{y, i} a_{i}=0  \tag{6}\\
& m \ddot{z}=\sum_{i=1}^{n} k_{z, i}\left(-z+b_{i} \alpha+a_{i} \beta\right)  \tag{7}\\
& m \ddot{z}+\sum_{i=1}^{n} k_{z, i}\left(z-b_{i} \alpha-a_{i} \beta\right)=0  \tag{8}\\
& m \ddot{z}+z \sum_{i=1}^{n} k_{z, i}-\alpha \sum_{i=1}^{n} k_{z, i} b_{i}-\beta \sum_{i=1}^{n} k_{z, i} a_{i}=0 \tag{9}
\end{align*}
$$

$$
\begin{align*}
& J_{x} \ddot{\alpha}=-\sum_{i=1}^{n} k_{y, i}\left(-y+c_{i} \alpha-a_{i} \theta\right) c_{i}-\sum_{i=1}^{n} k_{z, i}\left(-z+b_{i} \alpha+a_{i} \beta\right) b_{i}  \tag{10}\\
& J_{x} \ddot{\alpha}+\sum_{i=1}^{n} k_{y, i}\left(-y+c_{i} \alpha-a_{i} \theta\right) c_{i}+\sum_{i=1}^{n} k_{z, i}\left(-z+b_{i} \alpha+a_{i} \beta\right) b_{i}=0  \tag{11}\\
& J_{x} \ddot{\alpha}-y \sum_{i=1}^{n} k_{y, i} c_{i}-z \sum_{i=1}^{n} k_{z, i} b_{i} \\
& +\alpha \sum_{i=1}^{n}\left[\mathrm{k}_{\mathrm{y}, \mathrm{i}} \mathrm{c}_{\mathrm{i}}{ }^{2}+\mathrm{k}_{\mathrm{z}, \mathrm{i}} \mathrm{~b}_{\mathrm{i}}{ }^{2}\right]+\beta \sum_{\mathrm{i}=1}^{\mathrm{n}}\left[\mathrm{k}_{\mathrm{z}, \mathrm{i}} \mathrm{a}_{\mathrm{i}} \mathrm{~b}_{\mathrm{i}}\right]-\theta \sum_{\mathrm{i}=1}^{\mathrm{n}}\left[\mathrm{k}_{\mathrm{y}, \mathrm{i}} \mathrm{a}_{\mathrm{i}} \mathrm{c}_{\mathrm{i}}\right]=0  \tag{12}\\
& J_{y} \ddot{\beta}+\sum_{i=1}^{n} k_{x, i}\left(-x-c_{i} \beta-b_{i} \theta\right) c_{i}-\sum_{i=1}^{n} k_{z, i}\left(-z+b_{i} \alpha+a_{i} \beta\right) a_{i}  \tag{13}\\
& \mathrm{~J}_{\mathrm{y}} \ddot{\beta}+\mathrm{x} \sum_{\mathrm{i}=1}^{\mathrm{n}} \mathrm{k}_{\mathrm{x}, \mathrm{i}} \mathrm{c}_{\mathrm{i}}-\mathrm{z} \sum_{\mathrm{i}=1}^{\mathrm{n}} \mathrm{k}_{\mathrm{z}, \mathrm{i}} \mathrm{a}_{\mathrm{i}} \\
& +\alpha \sum_{i=1}^{n} \mathrm{k}_{\mathrm{z}, \mathrm{i}} \mathrm{a}_{\mathrm{i}} \mathrm{~b}_{\mathrm{i}}+\beta \sum_{\mathrm{i}=1}^{\mathrm{n}}\left[\mathrm{k}_{\mathrm{x}, \mathrm{i}} \mathrm{c}_{\mathrm{i}}{ }^{2}+\mathrm{k}_{\mathrm{z}, \mathrm{i}} \mathrm{a}_{\mathrm{i}}{ }^{2}\right]+\theta \sum_{\mathrm{i}=1}^{\mathrm{n}} \mathrm{k}_{\mathrm{x}, \mathrm{i}} \mathrm{~b}_{\mathrm{i}} \mathrm{c}_{\mathrm{i}}=0  \tag{14}\\
& J_{Z} \ddot{\theta}=\sum_{i=1}^{n} k_{x, i}\left(-x-c_{i} \beta-b_{i} \theta\right) b_{i}+\sum_{i=1}^{n} k_{y, i}\left(-y+c_{i} \alpha-a_{i} \theta\right) a_{i}  \tag{15}\\
& J_{Z} \ddot{\theta}+\sum_{i=1}^{n} k_{x, i}\left(x+c_{i} \beta+b_{i} \theta\right) b_{i}+\sum_{i=1}^{n} k_{y, i}\left(y-c_{i} \alpha+a_{i} \theta\right) a_{i}=0 \tag{16}
\end{align*}
$$

$$
\begin{align*}
& J_{z} \ddot{\theta}+x \sum_{i=1}^{n} k_{x, i} b_{i}+y \sum_{i=1}^{n} k_{y, i} a_{i} \\
& -\alpha \sum_{i=1}^{n} k_{y, i} a_{i} c_{i}+\beta \sum_{i=1}^{n} k_{x, i} b_{i} c_{i}+\theta \sum_{i=1}^{n}\left[k_{y, i} a_{i}^{2}+k_{x, i} b_{i}^{2}\right]=0 \tag{17}
\end{align*}
$$

Typically

$$
\begin{align*}
& \mathrm{k}_{\mathrm{x}, \mathrm{i}}=\mathrm{k}_{\mathrm{x}} \quad \text { for all } \mathrm{i}  \tag{18}\\
& \mathrm{k}_{\mathrm{y}, \mathrm{i}}=\mathrm{k}_{\mathrm{y}} \quad \text { for all } \mathrm{i}  \tag{19}\\
& \mathrm{k}_{\mathrm{z}, \mathrm{i}}=\mathrm{k}_{\mathrm{z}} \quad \text { for all } \mathrm{i} \tag{20}
\end{align*}
$$

By substitution, the equations of motion simplify to

$$
\begin{align*}
& m \ddot{x}+x n k_{x}+\beta k_{x} \sum_{i=1}^{n} c_{i}+\theta k_{x} \sum_{i=1}^{n} b_{i}=0  \tag{21}\\
& m \ddot{y}+y n k_{y}-\alpha k_{y} \sum_{i=1}^{n} c_{i}+\theta k_{y} \sum_{i=1}^{n} a_{i}=0  \tag{22}\\
& m \ddot{z}+\mathrm{znk}_{\mathrm{Z}}-\alpha k_{z} \sum_{i=1}^{n} b_{i}-\beta k_{z} \sum_{i=1}^{n} a_{i}=0 \tag{23}
\end{align*}
$$

$$
\mathrm{J}_{\mathrm{x}} \ddot{\alpha}-\mathrm{yk} \mathrm{k}_{\mathrm{y}} \sum_{\mathrm{i}=1}^{\mathrm{n}} \mathrm{c}_{\mathrm{i}}-\mathrm{zk} \mathrm{k}_{\mathrm{z}} \sum_{\mathrm{i}=1}^{\mathrm{n}} \mathrm{~b}_{\mathrm{i}}
$$

$$
\begin{equation*}
+\alpha \sum_{\mathrm{i}=1}^{\mathrm{n}}\left[\mathrm{k}_{\mathrm{y}} \mathrm{c}_{\mathrm{i}}^{2}+\mathrm{k}_{\mathrm{z}} \mathrm{~b}_{\mathrm{i}}^{2}\right]+\beta \mathrm{k}_{\mathrm{z}} \sum_{\mathrm{i}=1}^{\mathrm{n}}\left[\mathrm{a}_{\mathrm{i}} \mathrm{~b}_{\mathrm{i}}\right]-\theta \mathrm{k}_{\mathrm{y}} \sum_{\mathrm{i}=1}^{\mathrm{n}}\left[\mathrm{a}_{\mathrm{i}} \mathrm{c}_{\mathrm{i}}\right]=0 \tag{24}
\end{equation*}
$$

$$
\begin{align*}
& J_{y} \ddot{\beta}+x k_{x} \sum_{i=1}^{n} c_{i}-z_{k} \sum_{i=1}^{n} a_{i} \\
& +\alpha k_{z} \sum_{i=1}^{n} a_{i} b_{i}+\beta \sum_{i=1}^{n}\left[k_{X} c_{i}^{2}+k_{z} a_{i}^{2}\right]+\theta k_{x} \sum_{i=1}^{n} b_{i} c_{i}=0  \tag{25}\\
& J_{Z} \ddot{\theta}+x k_{x} \sum_{i=1}^{n} b_{i}+y k_{y} \sum_{i=1}^{n} a_{i} \\
& -\alpha k_{y} \sum_{i=1}^{n} a_{i} c_{i}+\beta k_{x} \sum_{i=1}^{n} b_{i} c_{i}+\theta \sum_{i=1}^{n}\left[k_{y} a_{i}^{2}+k_{x} b_{i}^{2}\right]=0 \tag{26}
\end{align*}
$$

The equations can be arranged in matrix format.

$$
\underline{\mathrm{M}}\left[\begin{array}{c}
\ddot{\mathrm{x}}  \tag{27}\\
\ddot{\mathrm{y}} \\
\ddot{\mathrm{z}} \\
\ddot{\alpha} \\
\ddot{\beta} \\
\ddot{\theta}
\end{array}\right]+\underline{\mathrm{K}}\left[\begin{array}{c}
\mathrm{x} \\
\mathrm{y} \\
\mathrm{z} \\
\alpha \\
\beta \\
\theta
\end{array}\right]=\left[\begin{array}{l}
0 \\
0 \\
0 \\
0 \\
0 \\
0
\end{array}\right]
$$

The mass and stiffness matrices are shown in upper triangular form due to symmetry.

$$
\underline{\mathbf{M}}=\left[\begin{array}{cccccc}
\mathrm{m} & 0 & 0 & 0 & 0 & 0  \tag{28}\\
& \mathrm{~m} & 0 & 0 & 0 & 0 \\
& & \mathrm{~m} & 0 & 0 & 0 \\
& & & \mathrm{~J}_{\mathrm{x}} & 0 & 0 \\
& & & & \mathrm{~J}_{\mathrm{y}} & 0 \\
& & & & & \mathrm{~J}_{\mathrm{z}}
\end{array}\right]
$$

## Base Excitation

Consider three separate base excitation cases, following the convention in Reference 1.
The mass and stiffness matrices in equations (28) and (29) apply in each of the cases.
Let $\mathrm{u}, \mathrm{x}$, and w be the base displacement in the $\mathrm{X}, \mathrm{Y}$ and Z -axes, respectively.

X-axis Excitation
Let

$$
\begin{equation*}
\mathrm{r}_{1}=\mathrm{x}-\mathrm{u} \tag{30}
\end{equation*}
$$

The equations can be arranged in matrix format.

$$
\underline{\mathrm{M}}\left[\begin{array}{c}
\ddot{\mathrm{r}}_{1}  \tag{31}\\
\ddot{\mathrm{y}} \\
\ddot{\mathrm{z}} \\
\ddot{\alpha} \\
\ddot{\beta} \\
\ddot{\theta}
\end{array}\right]+\underline{\mathrm{K}}\left[\begin{array}{c}
\mathrm{r}_{1} \\
\mathrm{y} \\
\mathrm{z} \\
\alpha \\
\beta \\
\theta
\end{array}\right]=\left[\begin{array}{c}
-\mathrm{m} \ddot{\mathrm{u}} \\
0 \\
0 \\
0 \\
0 \\
0
\end{array}\right]
$$

## Y-axis Excitation

Let

$$
\begin{equation*}
r_{2}=y-v \tag{32}
\end{equation*}
$$

The equations can be arranged in matrix format.

$$
\underline{\mathrm{M}}\left[\begin{array}{c}
\ddot{\mathrm{x}}  \tag{33}\\
\ddot{\mathrm{r}_{2}} \\
\ddot{\mathrm{z}} \\
\ddot{\alpha} \\
\ddot{\beta} \\
\ddot{\theta}
\end{array}\right]+\underline{\mathrm{K}}\left[\begin{array}{c}
\mathrm{x} \\
\mathrm{r}_{2} \\
\mathrm{z} \\
\alpha \\
\beta \\
\theta
\end{array}\right]=\left[\begin{array}{c}
0 \\
-\mathrm{m} \ddot{\mathrm{v}} \\
0 \\
0 \\
0 \\
0
\end{array}\right]
$$

Let

$$
\begin{equation*}
\mathrm{r}_{3}=\mathrm{z}-\mathrm{w} \tag{34}
\end{equation*}
$$

The equations can be arranged in matrix format.

$$
\underline{\mathrm{M}}\left[\begin{array}{c}
\ddot{\mathrm{x}}  \tag{35}\\
\ddot{\mathrm{y}} \\
\ddot{\mathrm{r}}_{3} \\
\ddot{\alpha} \\
\ddot{\beta} \\
\ddot{\theta}
\end{array}\right]+\underline{\mathrm{K}}\left[\begin{array}{c}
\mathrm{x} \\
\mathrm{y} \\
\mathrm{r}_{3} \\
\alpha \\
\beta \\
\theta
\end{array}\right]=\left[\begin{array}{c}
0 \\
0 \\
-\mathrm{m} \ddot{\mathrm{w}} \\
0 \\
0 \\
0
\end{array}\right]
$$

The acceleration transmissibility functions of equations (31), (33) and (35) can then be determined via Reference 2.

The modal transient response to a base input time history can be calculated via References 3 and 4.

## References

1. T. Irvine, Vibration Analysis of an Isolated Mass with Six Degrees of Freedom, Revision F, Vibrationdata, 2011.
2. T. Irvine, Frequency Response Function Analysis of a Multi-degree-of-freedom System with Enforced Motion, Vibrationdata, 2011.
3. T. Irvine, The Generalized Coordinate Method for Discrete Systems, Subjected to Base Excitation, Revision B, Vibrationdata, 2004.
4. T. Irvine, Shock Response of Multi-degree-of-freedom Systems, Revision F, Vibrationdata, 2010.

## APPENDIX A

## Example 1

A mass is mounted to a surface with six isolators. The system has the following properties.

| M | $=35 \mathrm{lbm}$ |
| ---: | :--- |
| Jx | $=804 \mathrm{lbm} \mathrm{in}^{\wedge} 2$ |
| Jy | $=1213 \mathrm{lbm} \mathrm{in}^{\wedge} 2$ |
| Jz | $=1035 \mathrm{lbm} \mathrm{in}^{\wedge} 2$ |
| kx | $=1200 \mathrm{lbf} / \mathrm{in}$ |
| ky | $=1600 \mathrm{lbf} / \mathrm{in}$ |
| kz | $=1200 \mathrm{lbf} / \mathrm{in}$ |

The six springs mount to the box at the following distances (inches) from the box C.G. The pattern is shown in Figure A-1.

| Isolator | X | Y | Z |
| :---: | :---: | :---: | :---: |
| 1 | 5.0 | -3.91 | -5.3 |
| 2 | 5.0 | -3.91 | 5.3 |
| 3 | 7.1 | -3.91 | 0 |
| 4 | -5.0 | -3.91 | 5.3 |
| 5 | -5.0 | -3.91 | -5.3 |
| 6 | -7.1 | -3.91 | 0 |

Assume uniform modal damping of $10 \%$, which is equivalent to $\mathrm{Q}=5$.
The results are calculated via Matlab script: arbitrary_isolators.m
A partial output listing is given on the next page. The results agree with a separate finite element model which is omitted for brevity.

ISOLATOR PATTERN EXAMPLE 1


Figure A-1.

The circles represent the isolator mounting points. The view is looking down along the Y -axis to the XZ plane.

The mass matrix is
$m=$

| 0.0907 | 0 | 0 | 0 | 0 | 0 |
| ---: | ---: | ---: | ---: | ---: | ---: |
| 0 | 0.0907 | 0 | 0 | 0 | 0 |
| 0 | 0 | 0.0907 | 0 | 0 | 0 |
| 0 | 0 | 0 | 2.0829 | 0 | 0 |
| 0 | 0 | 0 | 0 | 3.1425 | 0 |
| 0 | 0 | 0 | 0 | 0 | 2.6813 |

The stiffness matrix is
$\mathrm{k}=$
1.0e+005 *

| 0.0720 | 0 | 0 | 0 | 0 | 0.2815 |
| ---: | ---: | ---: | ---: | ---: | ---: |
| 0 | 0.0960 | 0 | 0 | 0 | 0.0000 |
| 0 | 0 | 0.0720 | -0.2815 | 0.0000 | 0 |
| 0 | 0 | -0.2815 | 2.8985 | 0.0000 | 0 |
| 0 | 0 | 0.0000 | 0.0000 | 3.7582 | 0 |
| 0.2815 | 0.0000 | 0 | 0 | 0 | 4.3139 |

Eigenvalues
lambda $=$
1.0e+005 *
$\begin{array}{llllll}0.3794 & 0.5001 & 1.0587 & 1.1959 & 1.8062 & 1.9028\end{array}$

Natural Frequencies $=$

1. 31 Hz
2. $\quad 35.59 \mathrm{~Hz}$
3. 51.79 Hz
4. $\quad 55.04 \mathrm{~Hz}$
5. $\quad 67.64 \mathrm{~Hz}$
6. $\quad 69.43 \mathrm{~Hz}$

| Modes Shapes | (rows represent modes) |  |  |  |  |
| ---: | ---: | ---: | ---: | ---: | ---: | ---: |
| x | y | z | alpha | beta | theta |
| 0 | 0 | 2.8 | 0.374 | 0 | 0 |
| -2.95 | 0 | 0 | 0 | 0 | 0.28 |
| 0 | -3.32 | 0 | 0 | 0 | 0 |
| 0 | 0 | 0 | 0 | -0.564 | 0 |
| 0 | 0 | -1.79 | 0.584 | 0 | 0 |
| -1.52 | 0 | 0 | 0 | 0 | -0.543 |



Figure A-2.

Both the Y and Z-axes responses were below the lower amplitude plot limit.


Figure A-3.

Both the X and Z -axes responses were below the lower amplitude plot limit.


Figure A-4.

Both the X and Y -axes responses were below the lower amplitude plot limit.

## APPENDIX B

## Example 2

A mass is mounted to a surface with six isolators. The system has the following properties.

| M | $=9.0 \mathrm{lbm}$ |
| :---: | :---: |
| JX | $=90.17 \mathrm{lbm}$ in^2 |
| Jy | $=105.2 \mathrm{lbm} \mathrm{in}$ ^2 |
| Jz | $=42.08 \mathrm{lbm}$ in^2 |
| kx | $=500 \mathrm{lbf} / \mathrm{in}$ |
| ky | $=500 \mathrm{lbf} / \mathrm{in}$ |
|  | $=500 \mathrm{lbf} / \mathrm{in}$ |

The six springs mount to the box at the following distances (inches) from the box C.G. The pattern is shown in Figure B-1.

| Isolator | X | Y | Z |
| :---: | :---: | :---: | :---: |
| 1 | -3 | -2 | -3 |
| 2 | -3 | -2 | 1 |
| 3 | -3 | -2 | 5 |
| 4 | 3 | -2 | -3 |
| 5 | 3 | -2 | 1 |
| 6 | 3 | -2 | 5 |

The results are calculated via Matlab script: arbitrary_isolators.m
A partial output listing is given on the next page. The results agree with the example in Reference 1.

ISOLATOR PATTERN EXAMPLE 2


Figure B-1.

The circles represent the isolator mounting points. The view is looking down along the Y -axis to the XZ plane.

The mass matrix is

```
m =
\begin{tabular}{rrrrrr}
0.0233 & 0 & 0 & 0 & 0 & 0 \\
0 & 0.0233 & 0 & 0 & 0 & 0 \\
0 & 0 & 0.0233 & 0 & 0 & 0 \\
0 & 0 & 0 & 0.2336 & 0 & 0 \\
0 & 0 & 0 & 0 & 0.2725 & 0 \\
0 & 0 & 0 & 0 & 0 & 0.109
\end{tabular}
The stiffness matrix is
\(\mathrm{k}=\)
\begin{tabular}{rrrrrr}
3000 & 0 & 0 & 0 & 3000 & 6000 \\
0 & 3000 & 0 & -3000 & 0 & 0 \\
0 & 0 & 3000 & -6000 & 0 & 0 \\
0 & -3000 & -6000 & 47000 & 0 & 0 \\
3000 & 0 & 0 & 0 & 62000 & 6000 \\
6000 & 0 & 0 & 0 & 6000 & 39000
\end{tabular}
Eigenvalues
lambda \(=\)
\(1.0 e+005\) *
\(0.6707 \quad 0.7501 \quad 2.2867 \quad 2.1907 \quad 2.6280 \quad 4.1983\)
Natural Frequencies \(=\)
1. 41.22 Hz
2. \(\quad 43.59 \mathrm{~Hz}\)
3. \(\quad 57.09 \mathrm{~Hz}\)
4. \(\quad 74.49 \mathrm{~Hz}\)
5. 81.59 Hz
6. 103.1 Hz
```

|  | X | Y | z | alpha | beta | theta |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 0 | 2.42 | 4.85 | 1.16 | 0 | 0 |
| 2 | 6.01 | 0 | 0 | 0 | -0.273 | -1.12 |
| 3 | 0 | -5.86 | 2.93 | 0 | 0 | 0 |
| 4 | -0.254 | 0 | 0 | 0 | -1.84 | 0.831 |
| 5 | 0 | -1.64 | -3.29 | 1.71 | 0 | 0 |
| 6 | 2.58 | 0 | 0 | 0 | 0.455 | 2.69 |

