

RING VIBRATION MODES
Revision D

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Variables

C	torsion constant
C_L	longitudinal wave speed
d	diameter
E	elastic modulus
G	shear modulus
f_n	natural frequency
f_r	ring frequency
h	height
I_x, I_y	cross-sectional area moment of inertia
m	mass per arc length
t	thickness
μ	mass per volume
ν	Poisson ratio

Ring Frequency

Consider a ring with a rectangular cross section and with completely free boundary conditions.

The ring frequency corresponds to the mode in which all points move radially outward together and then radially inward together. This is the first extension mode. It is analogous to a longitudinal mode in a rod.

The ring frequency f_r is the frequency at which the longitudinal wavelength in the skin material is equal to the vehicle circumference.

$$f_r = \frac{C_L}{\pi d} \quad (1)$$

Note that the wave speed can be calculated as

$$C_L = \sqrt{\frac{E}{\mu}} \quad (2)$$

The longitudinal wave speed in aluminum is approximately 16,700 feet per second.

Thus the ring frequency for aluminum is

$$f_r = \frac{5316 \text{ ft Hz}}{d} \quad (3)$$

$$f_r = \frac{63,790 \text{ in Hz}}{d} \quad (4)$$

Note that the ring frequency mode is a higher mode. It occurs at a much higher frequency than the first few bending modes.

Equations (1) through (4) are taken from Reference 1.

In-plane Bending Mode

The natural frequency equation for the in-plane bending modes is

$$f_n = \frac{2n(n^2 - 1)}{\pi d^2 \sqrt{n^2 + 1}} \sqrt{\frac{EI_y}{m}} \quad , \quad n = 1, 2, 3, \dots \quad (5)$$

Note that the Y-axis is the longitudinal axis, with its origin at the center of the cross-section.

Equation (5) is taken from Reference 2. Note that the $n=1$ case corresponds to in-plane rigid-body translation, with a frequency of zero.

The area moment of inertia for a thin ring is

$$I_y = \frac{1}{12} h t^3 \quad (6)$$

By substitution,

$$f_n = \frac{2n(n^2 - 1)}{\pi d^2 \sqrt{n^2 + 1}} \sqrt{\frac{E}{m} \left(\frac{1}{12} h t^3 \right)} \quad , \quad n = 1, 2, 3, \dots \quad (7)$$

Note that

$$m = \mu h t \quad (8)$$

By substitution,

$$f_n = \frac{2n(n^2 - 1)}{\pi d^2 \sqrt{n^2 + 1}} \sqrt{\mu h t \left(\frac{1}{12} h t^3 \right)} \quad , \quad n = 1, 2, 3, \dots \quad (9)$$

$$f_n = \frac{n(n^2 - 1)}{\pi d^2 \sqrt{n^2 + 1}} \sqrt{\frac{1}{3} \frac{E t^2}{\mu}} \quad , \quad n = 1, 2, 3, \dots \quad (10)$$

Out-of-Plane Bending Mode

The natural frequency equation for the out-of-plane bending modes is

$$f_n = \frac{2n(n^2 - 1)}{\pi d^2} \sqrt{\frac{EI_x}{m \left(n^2 + \frac{EI_x}{GC} \right)}} , \quad n = 1, 2, 3, \dots \quad (11)$$

Equation (11) is taken from Reference 2. Note that the $n=1$ case corresponds to out-of-plane rigid-body translation, with a frequency of zero.

For a rectangular cross section,

$$I_x = \frac{1}{12} t h^3 \quad (12)$$

$$f_n = \frac{2n(n^2 - 1)}{\pi d^2} \sqrt{\frac{\frac{1}{12} t h^3 E}{m \left(n^2 + \frac{\frac{1}{12} t h^3 E}{GC} \right)}} , \quad n = 1, 2, 3, \dots \quad (13)$$

$$f_n = \frac{2n(n^2 - 1)}{\pi d^2} \sqrt{\frac{t h^3 E}{m \left(12 n^2 + \frac{t h^3 E}{GC} \right)}} , \quad n = 1, 2, 3, \dots \quad (14)$$

$$f_n = \frac{2n(n^2 - 1)}{\pi d^2} \sqrt{\frac{t h^3 E}{\mu h t \left(12 n^2 + \frac{t h^3 E}{GC} \right)}}, \quad n = 1, 2, 3, \dots \quad (15)$$

$$f_n = \frac{2n(n^2 - 1)}{\pi} \left(\frac{h}{d^2} \right) \sqrt{\frac{E}{\mu \left(12 n^2 + \frac{t h^3 E}{GC} \right)}}, \quad n = 1, 2, 3, \dots \quad (16)$$

$$\frac{E}{G} = 2(1 + \nu) \quad (17)$$

By substitution,

$$f_n = \frac{2n(n^2 - 1)}{\pi} \left(\frac{h}{d^2} \right) \sqrt{\frac{E}{\mu \left(12 n^2 + \frac{2 t h^3 (1 + \nu)}{C} \right)}}, \quad n = 1, 2, 3, \dots \quad (18)$$

The torsion constant for a rectangular cross-section is taken as

$$C \approx \frac{0.3 t^3 h^3}{t^2 + h^2} \quad (19)$$

By substitution,

$$f_n = \frac{2n(n^2 - 1)}{\pi} \left(\frac{h}{d^2} \right) \sqrt{\frac{E}{\mu \left(12n^2 + \frac{2(t^2 + h^2)(1+\nu)}{0.3t^2} \right)}}, \quad n = 1, 2, 3, \dots$$

(20)

$$f_n = \frac{2n(n^2 - 1)}{\pi} \left(\frac{h}{d^2} \right) \sqrt{\frac{E}{\mu \left(12n^2 + 6.67 \left(1 + (h/t)^2 \right) (1+\nu) \right)}}, \quad n = 1, 2, 3, \dots$$

(21)

Model

A finite element was constructed as shown in Figure 1 and in Table 1.

Table 1. Model Parameters	
Parameter	Value
Number of Nodes	5040
Number of Elements	2016
Element Type	Solid
Thickness	0.25 inch
Height	1.0 inch
Diameter	40 inch
Boundary Condition	Completely Free
Material	Aluminum
Mass Density	0.098 lbm/in ³
Elastic Modulus	9900000. lbf/in ²
Software	FEMAP and NE/Nastran

The results of the finite element analysis are given in Figures 2 through 4.

The numerical results are summarized in Table 2, along with the theoretical values.

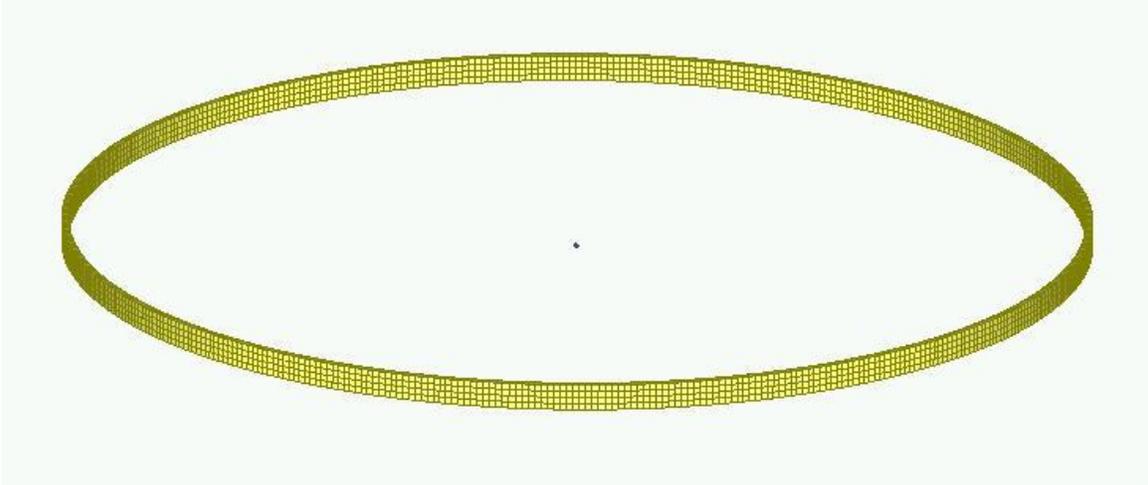


Figure 1. Undeformed Model

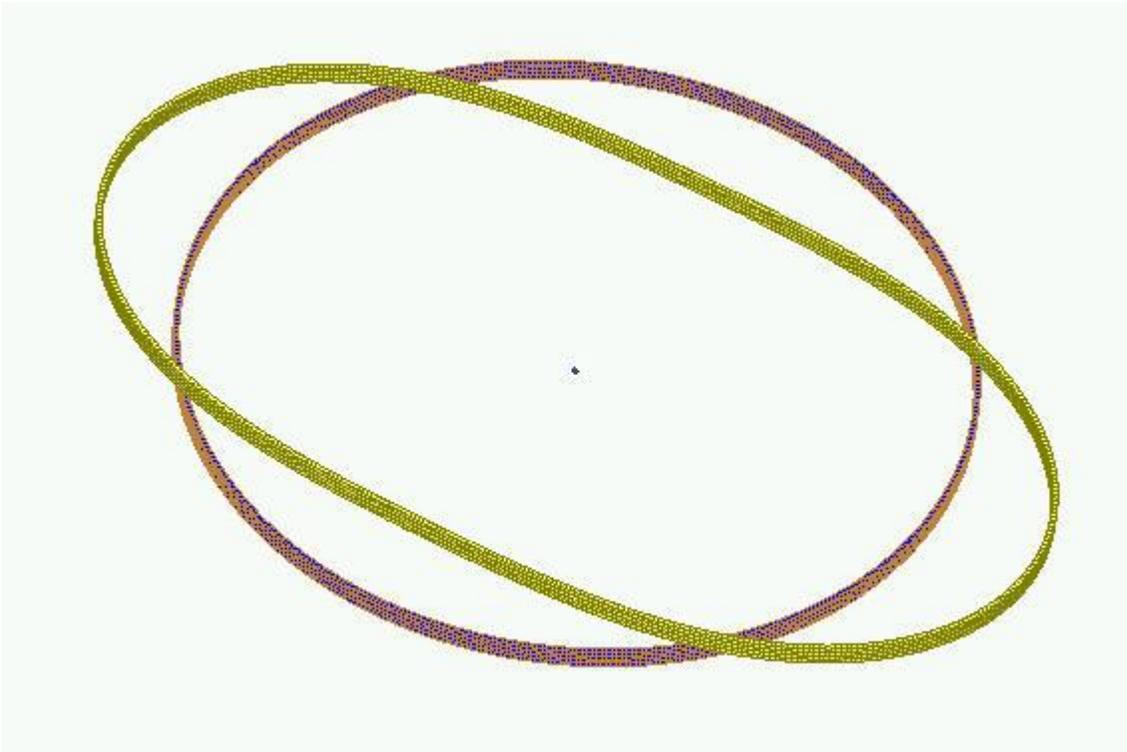


Figure 2. First In-plane Bending Mode, Superimposed on Undeformed Model, Frequency = 15.41 Hz

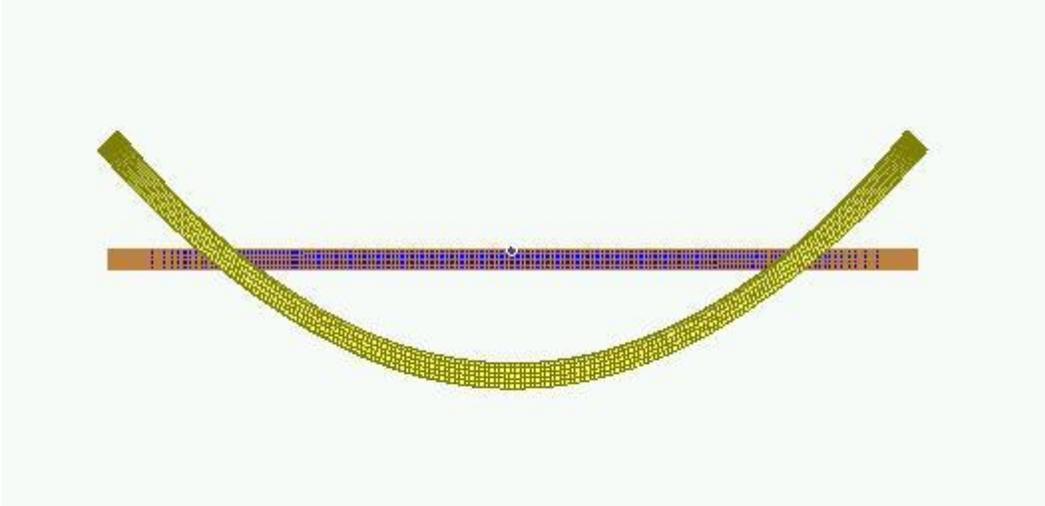


Figure 3. First Out-of-Plane Bending Mode, Superimposed on Undeformed Model, Frequency = 34.54 Hz

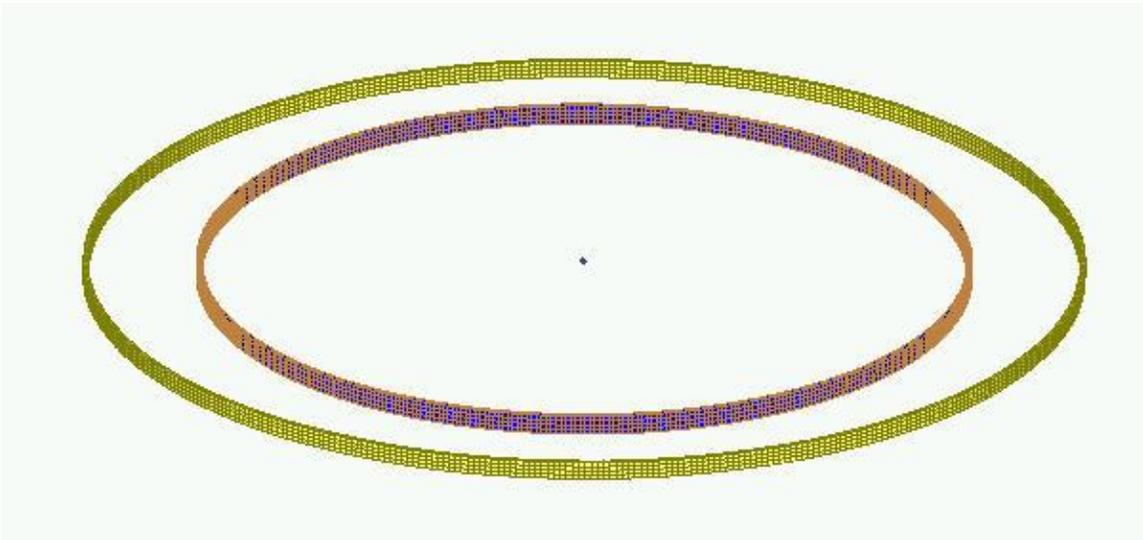


Figure 4. First Extensional Mode, Superimposed on Undeformed Model, Frequency = 1581 Hz

Table 2. Finite Element Model Results			
Mode	FEM Frequency (Hz)	Theoretical Frequency (Hz)	Theoretical Equation
First In-plane Bending	15.41	15.22	(5)
First Out-of-plane Bending	34.54	33.73	(11)
First Extensional (Ring Frequency)	1581.	1595.	(1)

The agreement is very good for each case.

References

1. T. Irvine, *Vibration Response of a Cylindrical Skin to Acoustic Pressure via the Franken Method*, Vibrationdata Publications, 2002.
2. R. Blevins, *Formulas for Natural Frequency and Mode Shape*, Krieger, Malabar, Florida, 1979.
3. Weaver, Timoshenko, and Young; *Vibration Problems in Engineering*, Wiley-Interscience, New York, 1990.

APPENDIX A

Extensional Vibration of a Ring

The derivation is taken from Reference 3.

u	is the radial displacement
E	is the modulus of elasticity
A	is the cross-sectional area
r	is the radius
d	is the diameter
ρ	is the mass/length
c	is the speed of sound in the material

The potential energy U is

$$U = \frac{AEu^2}{2r^2} 2\pi r \quad (\text{A-1})$$

$$U = \frac{\pi AE}{r} u^2 \quad (\text{A-2})$$

The kinetic energy T is

$$T = \frac{\rho A}{2} \dot{u}^2 2\pi r \quad (\text{A-3})$$

$$T = \pi r \rho A \dot{u}^2 \quad (\text{A-4})$$

Apply the energy method.

$$\frac{d}{dt}[\mathbf{T} + \mathbf{U}] = 0 \quad (\text{A-5})$$

$$\frac{d}{dt} \left[\pi r \rho A \dot{u}^2 + \frac{\pi A E}{r} u^2 \right] = 0 \quad (\text{A-6})$$

$$2\pi r \rho A \dot{u} \ddot{u} + 2 \frac{\pi A E}{r} u \dot{u} = 0 \quad (\text{A-7})$$

$$r \rho \ddot{u} + \frac{E}{r} u = 0 \quad (\text{A-8})$$

$$\ddot{u} + \frac{E}{\rho r^2} u = 0 \quad (\text{A-9})$$

$$\ddot{u} + \omega_n^2 u = 0 \quad (\text{A-10})$$

$$\omega_n^2 = \frac{E}{\rho r^2} \quad (\text{A-11})$$

$$\omega_n = \sqrt{\frac{E}{\rho r^2}} \quad (\text{A-12})$$

$$\omega_n = \frac{1}{r} \sqrt{\frac{E}{\rho}} \quad (\text{A-13})$$

$$f_n = \frac{1}{2\pi r} \sqrt{\frac{E}{\rho}} \quad (\text{A-14})$$

$$f_n = \frac{c}{\pi d} \quad (\text{A-15})$$

$$c = \sqrt{\frac{E}{\rho}} \quad (\text{A-16})$$

APPENDIX B

Fairing Ring Frequency Excited during Pyrotechnic Shock

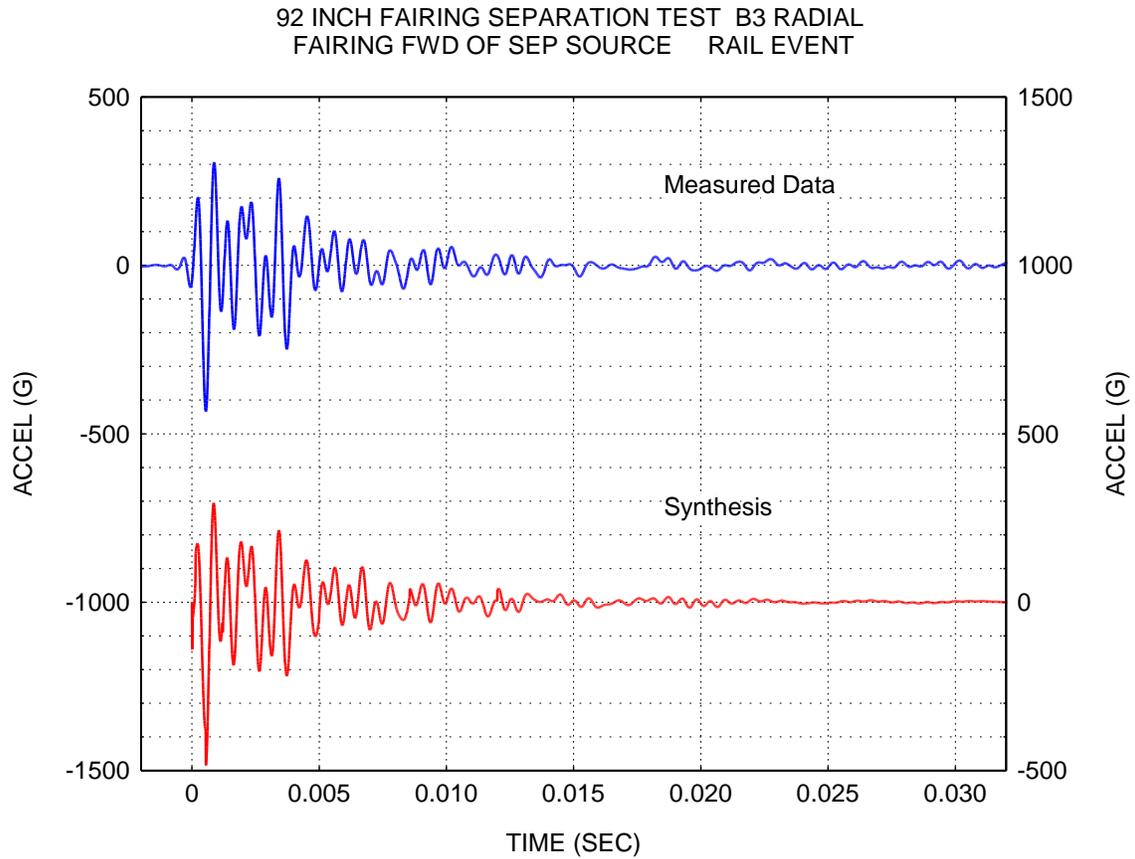


Figure B-1.

The source device was a frangible joint rail. The data was measured during a ground test.

The fairing consists of graphite-epoxy skins over an aluminum honeycomb core.

Note that the data is bandpass filtered from 20 to 2000 Hz.

The fairing's ring frequency is

$$f_n = \frac{C_L}{\pi d} \quad (B-1)$$

C_L	=	257,976 in/sec	Longitudinal wave speed in the composite skin material
d	=	92 in	Diameter

$$f_n = \frac{257,976 \text{ in/sec}}{\pi(92 \text{ in})} = 893 \text{ Hz}$$

The synthesis consists of ten components. The first three are given in Table B-1.

N	Amp (G)	Freq (Hz)	Phase (rad)	Damping	Delay (sec)
1	186.2	1888.895	1.682	0.017	0.000
2	172.7	897.063	3.609	0.037	0.000
3	112.6	1573.298	1.648	0.026	0.001

The second component represents the ring frequency.