

TWO-DEGREE-OF-FREEDOM SYSTEM
SUBJECTED TO A HALF-SINE PULSE FORCE
Revision A

By Tom Irvine
Email: tom@vibrationdata.com

March 21, 2014

Two-degree-of-freedom System

The method of generalized coordinates is demonstrated by an example. Consider the system in Figure 1.

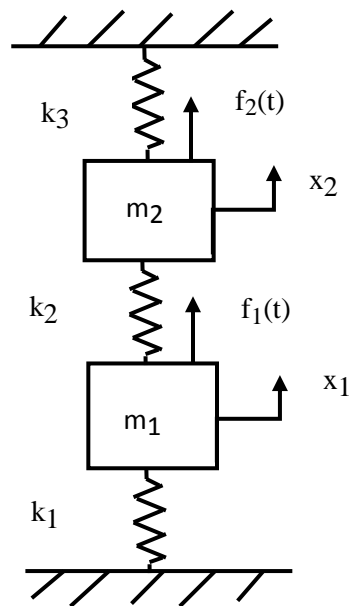


Figure 1.

A free-body diagram of mass 1 is given in Figure 2. A free-body diagram of mass 2 is given in Figure 3.

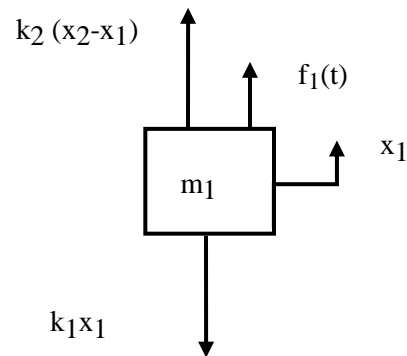


Figure 2.

Determine the equation of motion for mass 1.

$$\sum F = m_1 \ddot{x}_1 \quad (1)$$

$$m_1 \ddot{x}_1 = f_1(t) + k_2(x_2 - x_1) - k_1 x_1 \quad (2)$$

$$m_1 \ddot{x}_1 - k_2(x_2 - x_1) + k_1 x_1 = f_1(t) \quad (3)$$

$$m_1 \ddot{x}_1 + k_2(-x_2 + x_1) + k_1 x_1 = f_1(t) \quad (4)$$

$$m_1 \ddot{x}_1 + (k_1 + k_2)x_1 - k_2 x_2 = f_1(t) \quad (5)$$

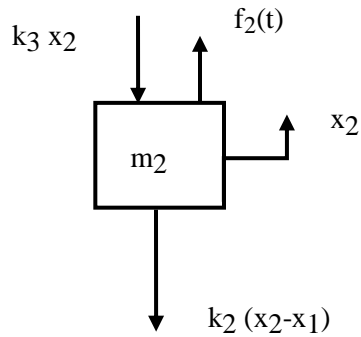


Figure 3.

Derive the equation of motion for mass 2.

$$\sum F = m_2 \ddot{x}_2 \quad (6)$$

$$m_2 \ddot{x}_2 = f_2(t) - k_2(x_2 - x_1) - k_3 x_2 \quad (7)$$

$$m_2 \ddot{x}_2 + k_2(x_2 - x_1) + k_3 x_2 = f_2(t) \quad (8)$$

$$m_2 \ddot{x}_2 + (k_2 + k_3)x_2 - k_2 x_1 = f_2(t) \quad (9)$$

Assemble the equations in matrix form.

$$\begin{bmatrix} m_1 & 0 \\ 0 & m_2 \end{bmatrix} \begin{bmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{bmatrix} + \begin{bmatrix} k_1 + k_2 & -k_2 \\ -k_2 & k_2 + k_3 \end{bmatrix} \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} = \begin{bmatrix} f_1(t) \\ f_2(t) \end{bmatrix} \quad (10)$$

Decoupling

Equation (10) is coupled via the stiffness matrix. An intermediate goal is to decouple the equation.

Simplify,

$$\mathbf{M}\ddot{\bar{x}} + \mathbf{K}\bar{x} = \bar{\mathbf{F}} \quad (11)$$

where

$$\mathbf{M} = \begin{bmatrix} m_1 & 0 \\ 0 & m_2 \end{bmatrix} \quad (12)$$

$$\mathbf{K} = \begin{bmatrix} k_1 + k_2 & -k_2 \\ -k_2 & k_2 + k_3 \end{bmatrix} \quad (13)$$

$$\bar{x} = \begin{bmatrix} x_1 \\ x_2 \end{bmatrix} \quad (14)$$

$$\bar{\mathbf{F}} = \begin{bmatrix} f_1(t) \\ f_2(t) \end{bmatrix} \quad (15)$$

Consider the homogeneous form of equation (11).

$$\mathbf{M}\ddot{\bar{x}} + \mathbf{K}\bar{x} = \bar{\mathbf{0}} \quad (16)$$

Seek a solution of the form

$$\bar{x} = \bar{q} \exp(j\omega t) \quad (17)$$

The q vector is the generalized coordinate vector.

Note that

$$\dot{\bar{x}} = j\omega \bar{q} \exp(j\omega t) \quad (18)$$

$$\ddot{\bar{x}} = -\omega^2 \bar{q} \exp(j\omega t) \quad (19)$$

Substitute equations (17) through (19) into equation (16).

$$-\omega^2 M \bar{q} \exp(j\omega t) + K \bar{q} \exp(j\omega t) = \bar{0} \quad (20)$$

$$\left\{ -\omega^2 M \bar{q} + K \bar{q} \right\} \exp(j\omega t) = \bar{0} \quad (21)$$

$$-\omega^2 M \bar{q} + K \bar{q} = \bar{0} \quad (22)$$

$$\left\{ -\omega^2 M + K \right\} \bar{q} = \bar{0} \quad (23)$$

$$\left\{ K - \omega^2 M \right\} \bar{q} = \bar{0} \quad (24)$$

Equation (24) is an example of a generalized eigenvalue problem. The eigenvalues can be found by setting the determinant equal to zero.

$$\det \left\{ K - \omega^2 M \right\} = \bar{0} \quad (25)$$

$$\det \left\{ \begin{bmatrix} k_1 + k_2 & -k_2 \\ -k_2 & k_2 + k_3 \end{bmatrix} - \omega^2 \begin{bmatrix} m_1 & 0 \\ 0 & m_2 \end{bmatrix} \right\} = 0 \quad (26)$$

$$\det \left\{ \begin{bmatrix} (k_1 + k_2) - \omega^2 m_1 & -k_2 \\ -k_2 & (k_2 + k_3) - \omega^2 m_2 \end{bmatrix} \right\} = 0 \quad (27)$$

$$\left[(k_1 + k_2) - \omega^2 m_1 \right] \left[(k_2 + k_3) - \omega^2 m_2 \right] - k_2^2 = 0 \quad (28)$$

$$(k_1 + k_2)(k_2 + k_3) - \omega^2 m_1(k_2 + k_3) - \omega^2 m_2(k_1 + k_2) + \omega^4 m_1 m_2 - k_2^2 = 0 \quad (29)$$

$$m_1 m_2 \omega^4 + [-m_1(k_2 + k_3) - m_2(k_1 + k_2)]\omega^2 + k_1 k_3 + (k_1 + k_3)k_2 + k_2^2 - k_2^2 = 0 \quad (30)$$

$$m_1 m_2 \omega^4 + [-m_1(k_2 + k_3) - m_2(k_1 + k_2)]\omega^2 + k_1 k_3 + k_1 k_2 + k_2 k_3 = 0 \quad (31)$$

The eigenvalues are the roots of the polynomial.

$$\omega_1^2 = \frac{-b - \sqrt{b^2 - 4ac}}{2a} \quad (32)$$

$$\omega_2^2 = \frac{-b + \sqrt{b^2 - 4ac}}{2a} \quad (33)$$

where

$$a = m_1 m_2 \quad (34)$$

$$b = [-m_1(k_2 + k_3) - m_2(k_1 + k_2)] \quad (35)$$

$$c = k_1 k_2 + k_1 k_3 + k_2 k_3 \quad (36)$$

The eigenvectors are found via the following equations.

$$\{K - \omega_1^2 M\} \bar{q}_1 = \bar{0} \quad (37)$$

$$\{K - \omega_2^2 M\} \bar{q}_2 = \bar{0} \quad (38)$$

where

$$\bar{q}_1 = \begin{bmatrix} v_1 \\ v_2 \end{bmatrix} \quad (39)$$

$$\bar{q}_2 = \begin{bmatrix} w_1 \\ w_2 \end{bmatrix} \quad (40)$$

An eigenvector matrix Q can be formed. The eigenvectors are inserted in column format.

$$Q = [\bar{q}_1 \quad \bar{q}_2] \quad (41)$$

$$Q = \begin{bmatrix} v_1 & w_1 \\ v_2 & w_2 \end{bmatrix} \quad (42)$$

The eigenvectors represent orthogonal mode shapes.

Each eigenvector can be multiplied by an arbitrary scale factor. A mass-normalized eigenvector matrix \hat{Q} can be obtained such that the following orthogonality relations are obtained.

$$\hat{Q}^T M \hat{Q} = I \quad (43)$$

and

$$\hat{Q}^T K \hat{Q} = \Omega \quad (44)$$

where

- I is the identity matrix
- Ω is a diagonal matrix of eigenvalues

The superscript T represents transpose.

Note the mass-normalized forms

$$\hat{Q} = \begin{bmatrix} \hat{v}_1 & \hat{w}_1 \\ \hat{v}_2 & \hat{w}_2 \end{bmatrix} \quad (45)$$

$$\hat{Q}^T = \begin{bmatrix} \hat{v}_1 & \hat{v}_2 \\ \hat{w}_1 & \hat{w}_2 \end{bmatrix} \quad (46)$$

Rigorous proof of the orthogonality relationships is beyond the scope of this tutorial.

Further discussion is given in References 1 and 2.

Nevertheless, the orthogonality relationships are demonstrated by an example in this tutorial.

Now define a generalized coordinate $\eta(t)$ such that

$$\bar{x} = \hat{Q} \bar{\eta} \quad (47)$$

Substitute equation (47) into the equation of motion, equation (11).

$$M \hat{Q} \bar{\ddot{\eta}} + K \hat{Q} \bar{\eta} = \bar{F} \quad (48)$$

Premultiply by the transpose of the normalized eigenvector matrix.

$$\hat{Q}^T M \hat{Q} \bar{\ddot{\eta}} + \hat{Q}^T K \hat{Q} \bar{\eta} = \hat{Q}^T \bar{F} \quad (49)$$

The orthogonality relationships yield

$$I \bar{\ddot{\eta}} + \Omega \bar{\eta} = \hat{Q}^T \bar{F} \quad (50)$$

The equations of motion along with an added damping matrix become

$$\begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} \ddot{\eta}_1 \\ \ddot{\eta}_2 \end{bmatrix} + \begin{bmatrix} 2\xi_1\omega_1 & 0 \\ 0 & 2\xi_2\omega_2 \end{bmatrix} \begin{bmatrix} \dot{\eta}_1 \\ \dot{\eta}_2 \end{bmatrix} + \begin{bmatrix} \omega_1^2 & 0 \\ 0 & \omega_2^2 \end{bmatrix} \begin{bmatrix} \eta_1 \\ \eta_2 \end{bmatrix} = \begin{bmatrix} \hat{v}_1 & \hat{v}_2 \\ \hat{w}_1 & \hat{w}_2 \end{bmatrix} \begin{bmatrix} f_1(t) \\ f_2(t) \end{bmatrix} \quad (51)$$

Note that the two equations are decoupled in terms of the generalized coordinate.

Equation (51) yields two equations

$$\ddot{\eta}_1 + 2\xi_1\omega_1\dot{\eta}_1 + \omega_1^2\eta_1 = \hat{v}_1f_1(t) + \hat{v}_2f_2(t) \quad (52)$$

$$\ddot{\eta}_2 + 2\xi_2\omega_2\dot{\eta}_2 + \omega_2^2\eta_2 = \hat{w}_1f_1(t) + \hat{w}_2f_2(t) \quad (53)$$

The equations can be solved in terms of Laplace transforms, or some other differential equation solution method.

Now consider the initial conditions. Recall

$$\bar{x} = \hat{Q} \bar{\eta} \quad (54)$$

Thus

$$\bar{x}(0) = \hat{Q} \bar{\eta}(0) \quad (55)$$

Premultiply by $\hat{Q}^T M$.

$$\hat{Q}^T M \bar{x}(0) = \hat{Q}^T M \hat{Q} \bar{\eta}(0) \quad (56)$$

Recall

$$\hat{Q}^T M \hat{Q} = I \quad (57)$$

$$\hat{Q}^T M \bar{x}(0) = I \bar{\eta}(0) \quad (58)$$

$$\hat{Q}^T M \bar{x}(0) = \bar{\eta}(0) \quad (59)$$

Finally, the transformed initial displacement is

$$\bar{\eta}(0) = \hat{Q}^T M \bar{x}(0) \quad (60)$$

Similarly, the transformed initial velocity is

$$\dot{\bar{\eta}}(0) = \hat{Q}^T M \dot{\bar{x}}(0) \quad (61)$$

A basis for a solution is thus derived.

Half-Sine Force

Now consider the special case of a half-sine force applied to mass 2.

$$f_1(t) = 0 \quad (62)$$

$$f_2(t) = B_2 \sin(\beta t), \quad \text{for } t \leq T \quad (63)$$

where

$$\beta = \pi/T$$

Thus,

$$\ddot{\eta}_1 + 2\xi_1 \omega_1 \dot{\eta}_1 + \omega_1^2 \eta_1 = \hat{v}_2 B_2 \sin(\beta t) \quad (64)$$

$$\ddot{\eta}_2 + 2\xi_2 \omega_2 \dot{\eta}_2 + \omega_2^2 \eta_2 = \hat{w}_2 B_2 \sin(\beta t) \quad (65)$$

Take the Laplace transform of equation (64).

$$\ddot{\eta}_1 + 2\xi_1 \omega_1 \dot{\eta}_1 + \omega_1^2 \eta_1 = \hat{v}_2 B_2 \sin(\beta t) \quad (66)$$

$$L\{\ddot{\eta}_1 + 2\xi_1\omega_1\dot{\eta}_1 + \omega_1^2\eta_1\} = L\{\hat{v}_2B_2\sin(\beta t)\} \quad (67)$$

$$\begin{aligned} & s^2\hat{\eta}_1(s) - s\eta_1(0) - \dot{\eta}_1(0) \\ & + 2\xi_1\omega_1s\hat{\eta}_1(s) - 2\xi_1\omega_1\eta_1(0) \\ & + \omega_1^2\hat{\eta}_1(s) = \hat{v}_2B_2\left\{\frac{\beta}{s^2 + \beta^2}\right\} \end{aligned} \quad (68)$$

$$\left\{s^2 + 2\xi_1\omega_1s + \omega_n^2\right\}\hat{\eta}_1(s) - \{s + 2\xi_1\omega_1\}\eta_1(0) - \dot{\eta}_1(0) = \hat{v}_2B_2\left\{\frac{\beta}{s^2 + \beta^2}\right\} \quad (69)$$

$$\left\{s^2 + 2\xi_1\omega_1s + \omega_n^2\right\}\hat{\eta}_1(s) = \hat{v}_2B_2\left\{\frac{\beta}{s^2 + \beta^2}\right\} + \{s + 2\xi_1\omega_1\}\eta_1(0) + \dot{\eta}_1(0) \quad (70)$$

$$\hat{\eta}_1(s) = \frac{\hat{v}_2B_2}{\left\{s^2 + 2\xi_1\omega_1s + \omega_n^2\right\}}\left\{\frac{\beta}{s^2 + \beta^2}\right\} + \frac{\{s + 2\xi_1\omega_1\}\eta_1(0) + \dot{\eta}_1(0)}{\left\{s^2 + 2\xi_1\omega_1s + \omega_n^2\right\}} \quad (71)$$

The solution is found via References 3 and 4. The inverse Laplace transform for the first modal coordinate is

$$\begin{aligned}
 \eta_1(t) = & \\
 & + \frac{\hat{v}_2 B_2}{\left[(\beta^2 - \omega_1^2)^2 + (2\xi_1 \beta \omega_1)^2 \right]} \left\{ -[2\xi_1 \omega_1 \beta] \cos(\beta t) - [\beta^2 - \omega_1^2] \sin(\beta t) \right\} \\
 & + \frac{\hat{v}_2 B_2 \left[\frac{\beta}{\omega_{d,1}} \right] \exp(-\xi_1 \omega_1 t)}{\left[(\beta^2 - \omega_1^2)^2 + (2\xi_1 \beta \omega_1)^2 \right]} \left\{ [2\xi_1 \omega_1 \omega_{d,1}] \cos(\omega_{d,1} t) + [\beta^2 - \omega_1^2 (1 - 2\xi_1^2)] \sin(\omega_{d,1} t) \right\} \\
 & + \exp(-\xi_1 \omega_1 t) \left\{ \eta_1(0) \cos(\omega_{d,1} t) + \left\{ \frac{\dot{\eta}_1(0) + (\xi_1 \omega_1) \eta_1(0)}{\omega_{d,1}} \right\} \sin(\omega_{d,1} t) \right\}, \\
 & \text{for } t \leq T
 \end{aligned}$$

(72)

Similarly,

$$\begin{aligned}
\eta_2(t) = & \\
& + \frac{\hat{w}_2 B_2}{\left[(\beta^2 - \omega_2^2)^2 + (2\xi_2 \beta \omega_2)^2 \right]} \left\{ -[2\xi_2 \omega_2 \beta] \cos(\beta t) - \frac{1}{\beta} [\beta^2 - \omega_2^2] \sin(\beta t) \right\} \\
& + \frac{\hat{w}_2 B_2 \left[\frac{\beta}{\omega_{d,2}} \right] \exp(-\xi_2 \omega_2 t)}{\left[(\beta^2 - \omega_2^2)^2 + (2\xi_2 \beta \omega_2)^2 \right]} \left\{ [2\xi_2 \omega_2 \omega_{d,2}] \cos(\omega_{d,2} t) + [\beta^2 - \omega_2^2 (1 - 2\xi_2^2)] \sin(\omega_{d,2} t) \right\} \\
& + \exp(-\xi_2 \omega_2 t) \left\{ \eta_2(0) \cos(\omega_{d,2} t) + \left\{ \frac{\dot{\eta}_2(0) + (\xi_2 \omega_2) \eta_2(0)}{\omega_{d,2}} \right\} \sin(\omega_{d,2} t) \right\}
\end{aligned}$$

for $t \leq T$

(73)

The physical displacements are found via

$$\bar{x} = \hat{Q} \bar{\eta} \quad (74)$$

Free Vibration

For $t > T$ and $\tau = t - T$

$$\eta_1(t) = \exp(-\xi_1 \omega_1 \tau) \left\{ \eta_1(T) \cos(\omega_{d,1} \tau) + \left\{ \frac{\dot{\eta}_1(T) + (\xi_1 \omega_1) \eta_1(T)}{\omega_{d,1}} \right\} \sin(\omega_{d,1} \tau) \right\} \quad (75)$$

$$\eta_2(t) = \exp(-\xi_2 \omega_2 \tau) \left\{ \eta_2(T) \cos(\omega_{d,2} \tau) + \left\{ \frac{\dot{\eta}_2(T) + (\xi_2 \omega_2) \eta_2(T)}{\omega_{d,2}} \right\} \sin(\omega_{d,2} \tau) \right\} \quad (76)$$

References

1. Bathe, Finite Element Procedures in Engineering Analysis, Prentice-Hall, New Jersey, 1982. Section 12.3.1.
2. Weaver and Johnston, Structural Dynamics by Finite Elements, Prentice-Hall, New Jersey, 1987. Chapter 4.
3. T. Irvine, Table of Laplace Transforms, Vibrationdata, 2000.
4. T. Irvine, Partial Fraction Expansion, Rev F, Vibrationdata, 2010.

APPENDIX A

Example

Consider the system in Figure 1 with the values in Table A-1.

Assume 5% damping for each mode. Assume zero initial conditions.

Table A-1. Parameters		
Variable	Value	Unit
m_1	3.0	lbf sec ² /in
m_2	2.0	lbf sec ² /in
k_1	400,000	lbf/in
k_2	300,000	lbf/in
k_3	100,000	lbf/in
B_2	100	lbf
T	0.011	sec

The mass matrix is

$$M = \begin{bmatrix} m_1 & 0 \\ 0 & m_2 \end{bmatrix} = \begin{bmatrix} 3 & 0 \\ 0 & 2 \end{bmatrix} \quad (\text{A-1})$$

The stiffness matrix is

$$K = \begin{bmatrix} k_1 + k_2 & -k_2 \\ -k_2 & k_2 + k_3 \end{bmatrix} = \begin{bmatrix} 700,000 & -300,000 \\ -300,000 & 400,000 \end{bmatrix} \quad (\text{A-2})$$

The analysis is performed using a Matlab script.

```
>> twodof_half_sine_force
```

```
twodof_half_sine_force.m  ver 1.0  January 18, 2012  
by Tom Irvine  Email: tomirvine@aol.com
```

```
This script calculates the response of a two-degree-of-freedom  
system to a half-sine force excitation at dof 2.
```

```
Enter the units system
```

```
1=English  2=metric
```

```
1
```

```
Assume symmetric mass and stiffness matrices.
```

```
Select input mass unit
```

```
1=lbm  2=lbm sec^2/in
```

```
2
```

```
stiffness unit = lbf/in
```

```
Select file input method
```

```
1=file preloaded into Matlab
```

```
2=Excel file
```

```
1
```

```
Mass Matrix
```

```
Enter the matrix name:  m_two
```

```
Stiffness Matrix
```

```
Enter the matrix name:  k_two
```

```
Input Matrices
```

```
mass =
```

```
    3    0  
    0    2
```

```
stiff =
```

```
    700000   -300000  
   -300000    400000
```

```
Natural Frequencies
```

```
No.      f (Hz)  
1.       48.552  
2.       92.839
```

```
Modes Shapes (column format)
```

```
ModeShapes =
```

```
    0.3797   -0.4349  
    0.5326    0.4651
```


Enter the damping ratio for mode 1 0.05
Enter the damping ratio for mode 2 0.05

Participation Factors =

2.204
-0.3746

Enter the force amplitude (lbf) 100
Enter the half-sine pulse duration (sec) 0.011

Enter the first initial displacement (in) 0
Enter the second initial displacement (in) 0

Enter the first initial velocity (in/sec) 0
Enter the second initial velocity (in/sec) 0

Enter the sample rate (samples/sec) 10000

Enter the analysis duration (sec) 0.15

dof 1 displacement (in)
max= 0.0003287
min= -0.0003149

dof 2 displacement (in)
max= 0.0005005
min= -0.0003728

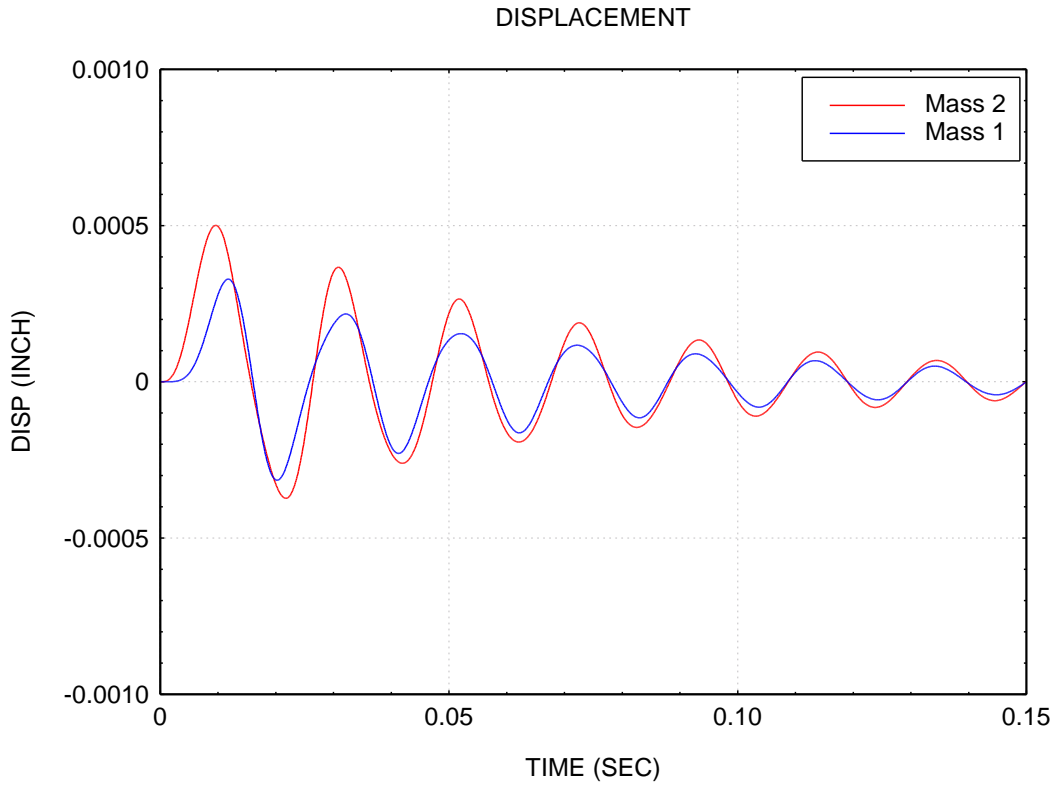


Figure A-1.

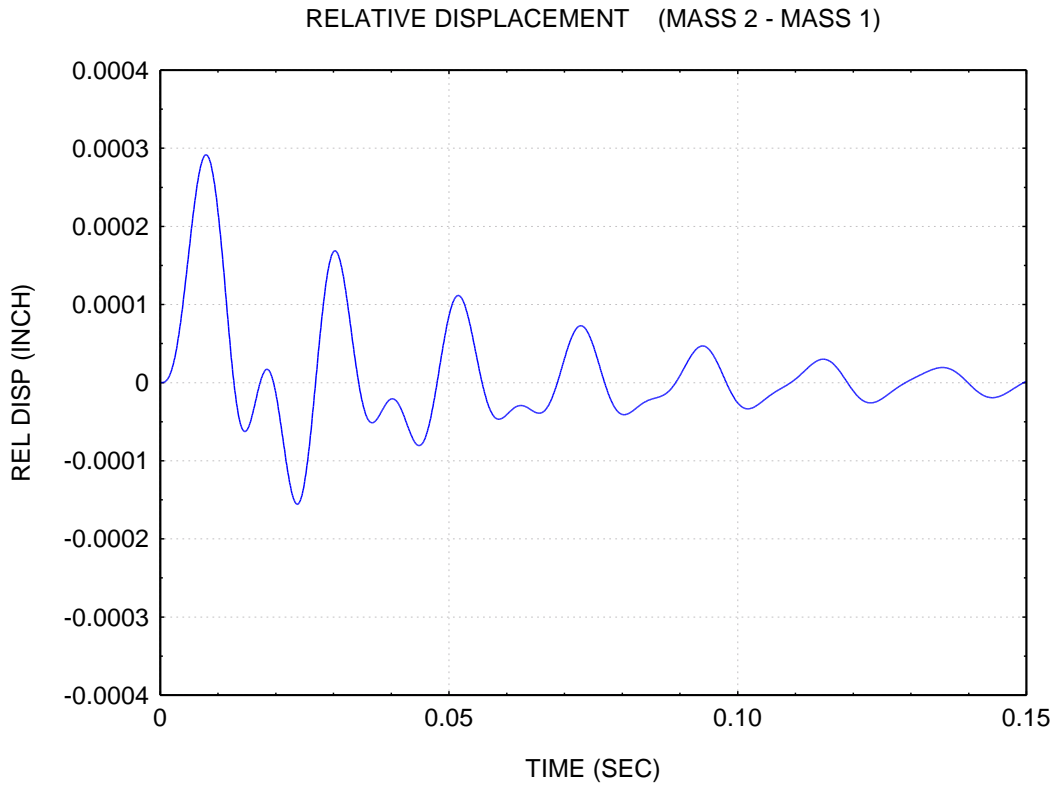


Figure A-2.

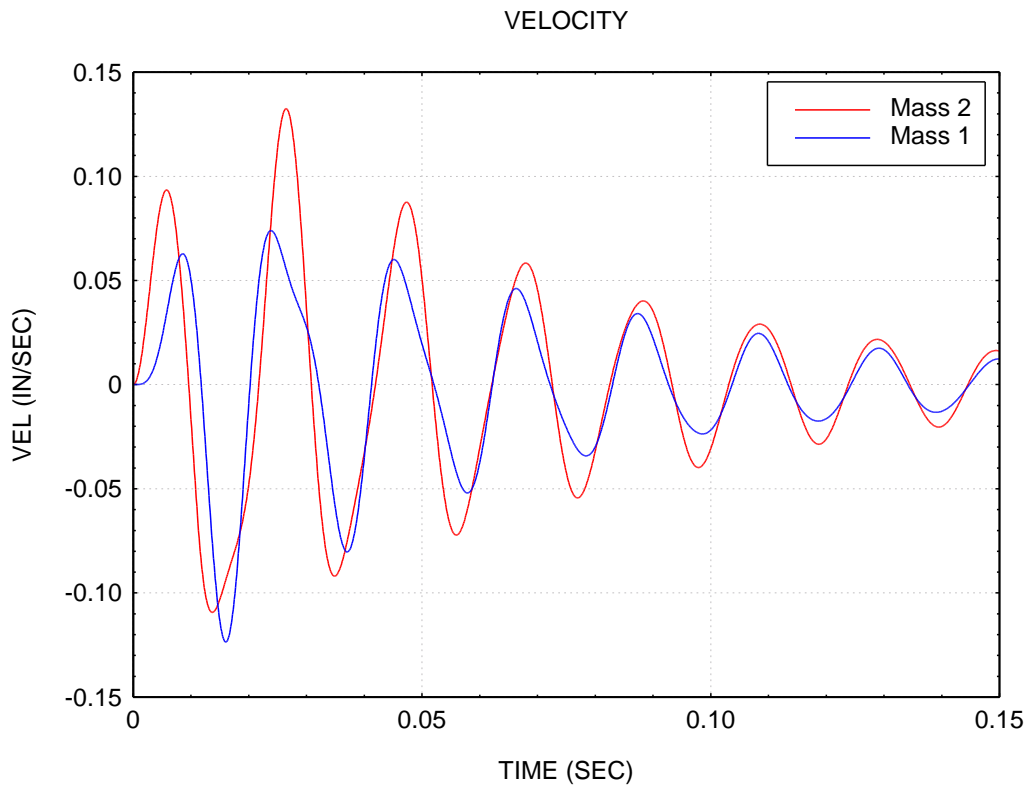


Figure A-3.

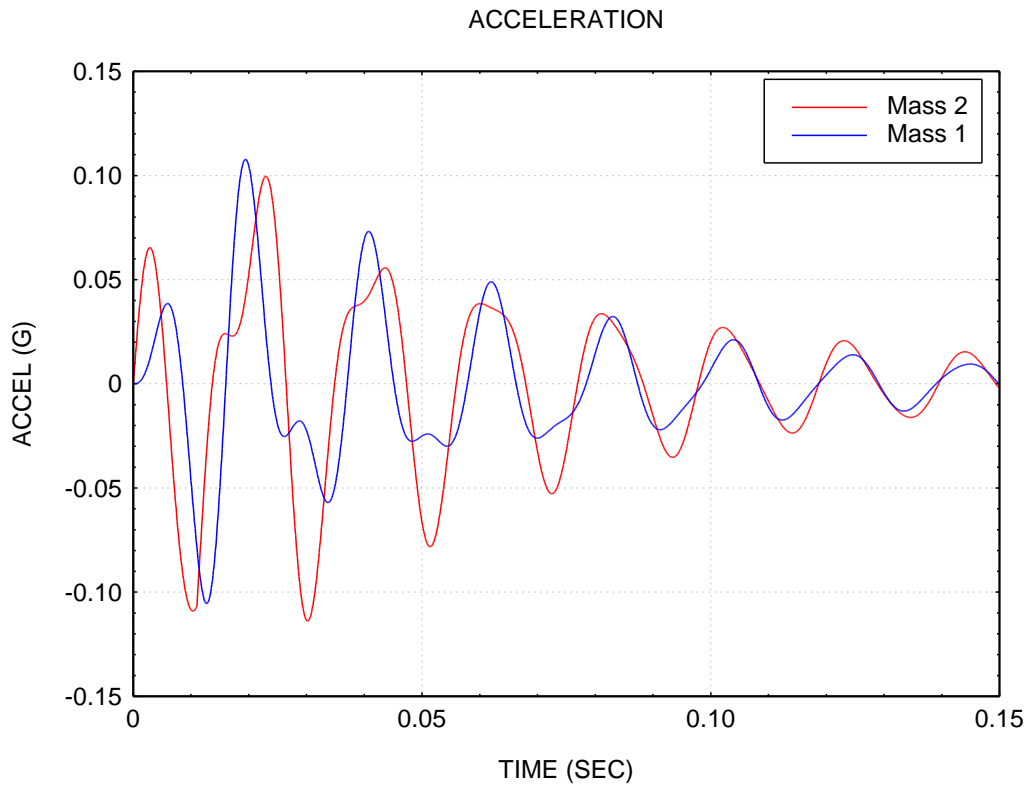


Figure A-4.